



Universidad
Católica del Norte



lawphysics
Latin American Webinars on Physics

Astroparticle physics

Roberto A. Lineros

Departamento de Física, Universidad Católica del Norte

Winter School HEP-PUC 2019 “Topics on Graviticulas”



The Plan

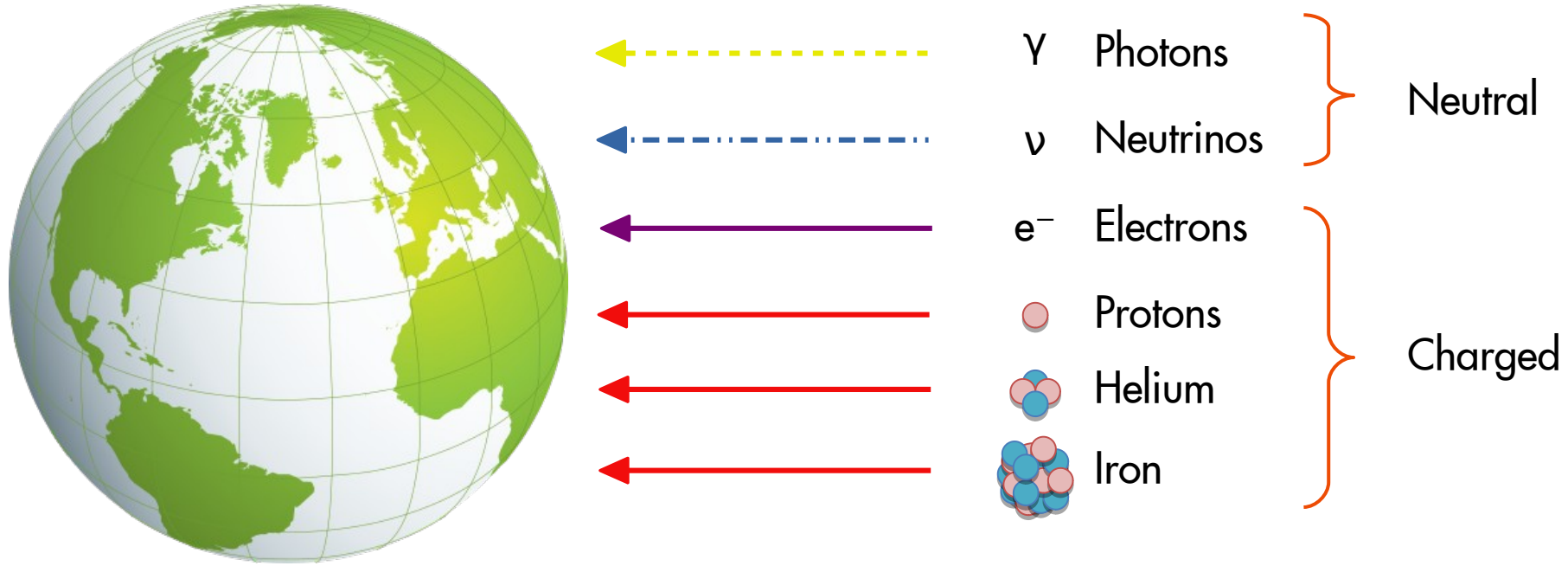
1. Astroparticles
2. Cosmic-rays
3. Neutrinos
4. Gamma-rays
5. Gravitational Waves

Dark Matter
Astrophysical Sources

Neutrinos

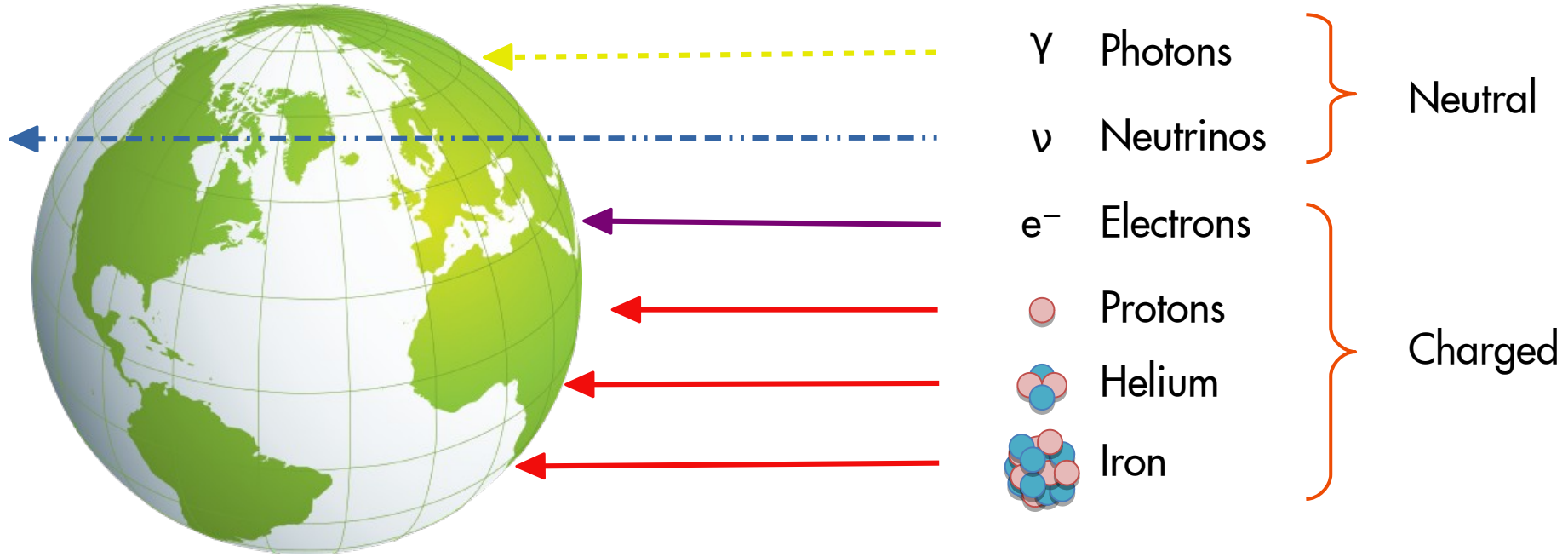


Particles from outer space



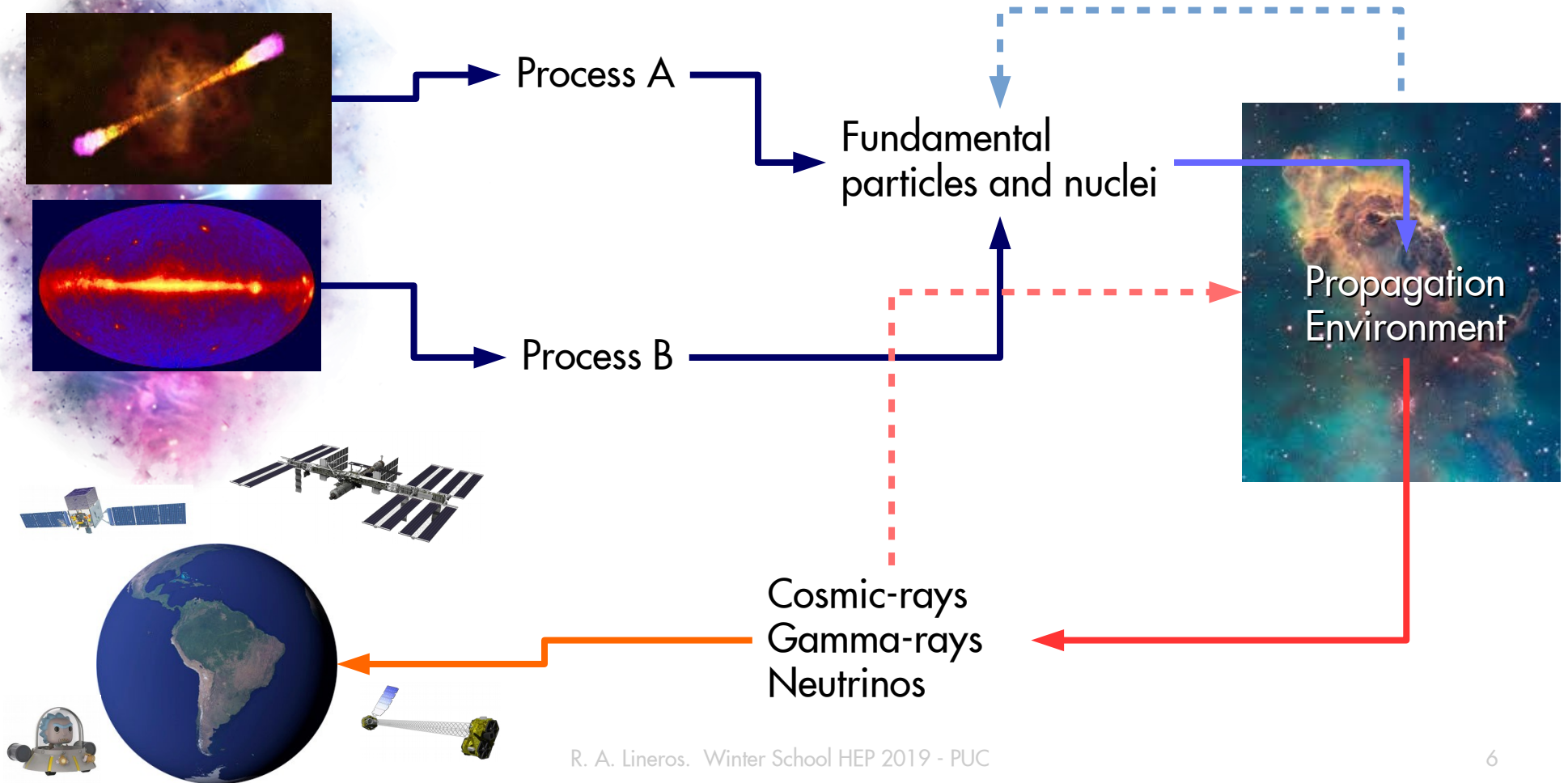
Astroparticles = Multimessenger and Multiwavelength

Particles from outer space



Neutrino can pass through Earth without any interaction

Multimessengers

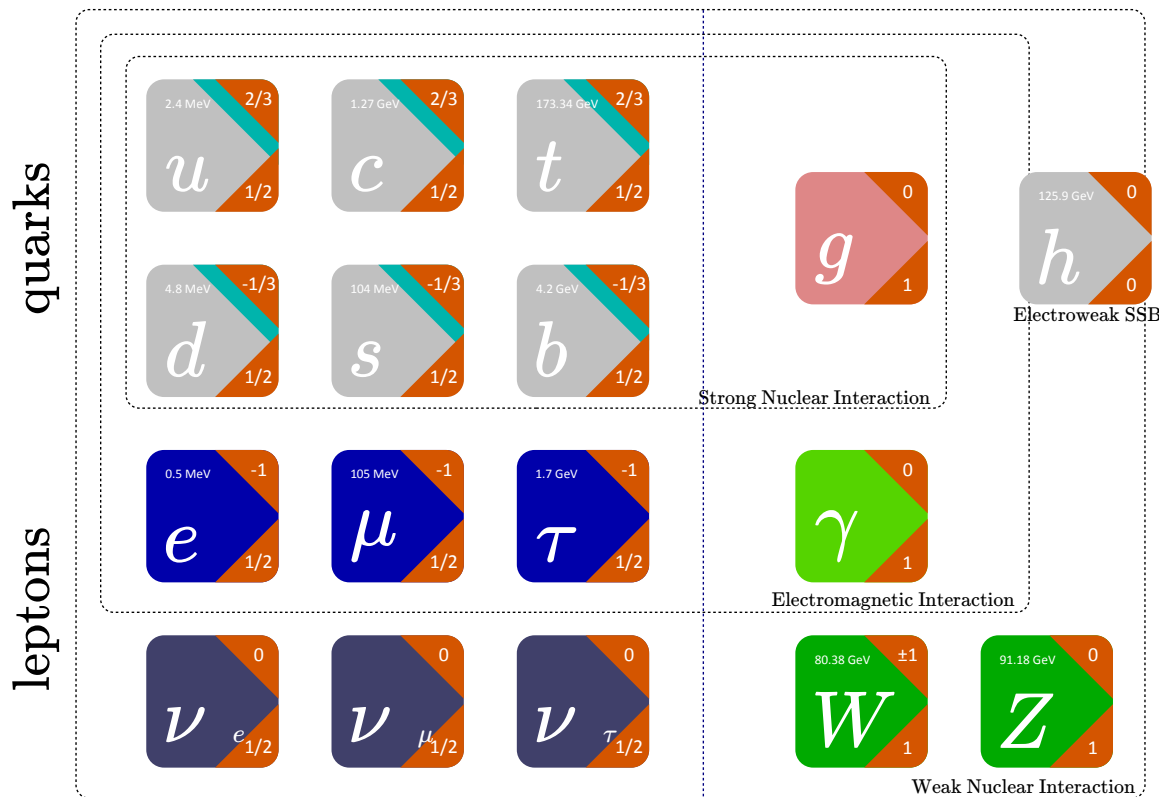


Neutrino Physics



The Standard Model

SM matter families



Symmetries

- CPT
- $SU(3)_c$: Color
- $SU(2)_L$: Isospin
- $U(1)_Y$: Hypercharge

Matter content

- 3 families quarks
- 3 families leptons

Higgs fields

- $SU(2)_L \times U(1)_Y \rightarrow U(1)_{EM}$
- Mass to fundamental particles

Neutrinos in the SM

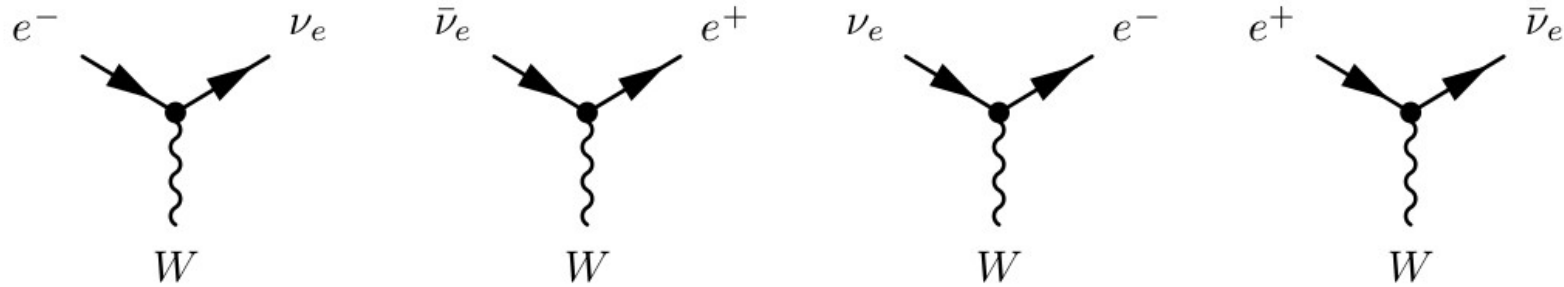
Neutrinos are part of the SU(2) of left-handed leptons

Fermion IRs under $SU(2)_L \times U(1)_Y$			SU(2)	isospin	hypercharge	electric charge
			I	I_3	Y	Q
$L_{eL} \equiv \begin{pmatrix} \nu_{eL} \\ e_L \end{pmatrix}$	$L_{\mu L} \equiv \begin{pmatrix} \nu_{\mu L} \\ \mu_L \end{pmatrix}$	$L_{\tau L} \equiv \begin{pmatrix} \nu_{\tau L} \\ \tau_L \end{pmatrix}$	1/2	1/2 -1/2	-1	0 -1
$l_{eR} \equiv e_R$	$l_{\mu R} \equiv \mu_R$	$l_{\tau R} \equiv \tau_R$	0	0	-2	-1

$$\mathcal{L} \propto \underbrace{i\bar{L}_\alpha \gamma_\mu D^\mu L_\alpha + i\bar{l}_{R\alpha} \gamma_\mu D^\mu l_{R\alpha}}_{\text{Covariant derivatives}} - \underbrace{(y_{\alpha\beta} \bar{L}_\alpha \Phi l_{R\beta} + h.c.)}_{\text{Interaction/mass terms}}$$

Neutrinos in the SM

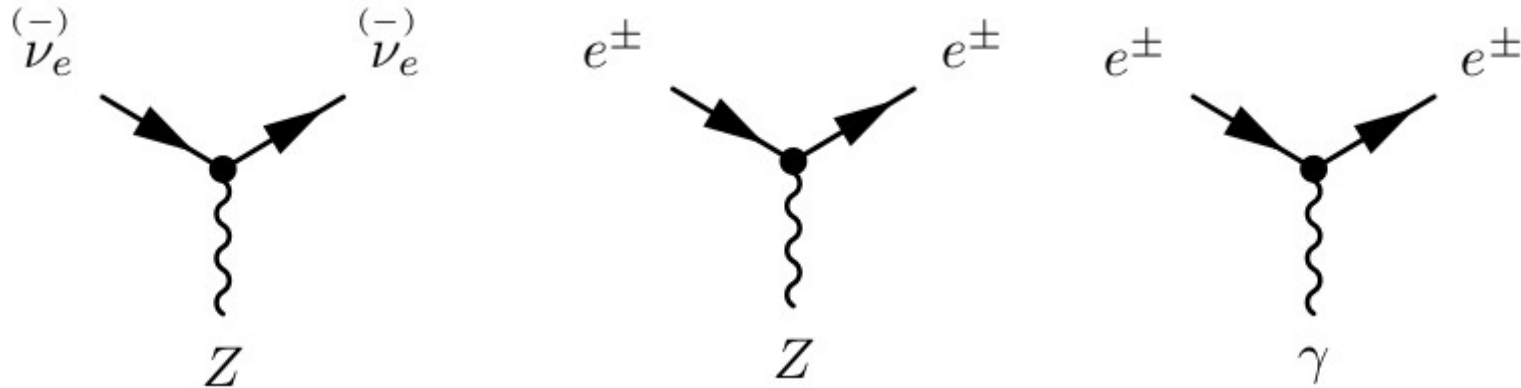
Standard Model Neutrinos only interact via weak boson: Z and W



$$\mathcal{L}_{I,L}^{(CC)} = -\frac{g}{\sqrt{2}} \{ \bar{\nu}_{eL} W e_L + \bar{e}_L W^\dagger \nu_{eL} \}$$

Neutrinos in the SM

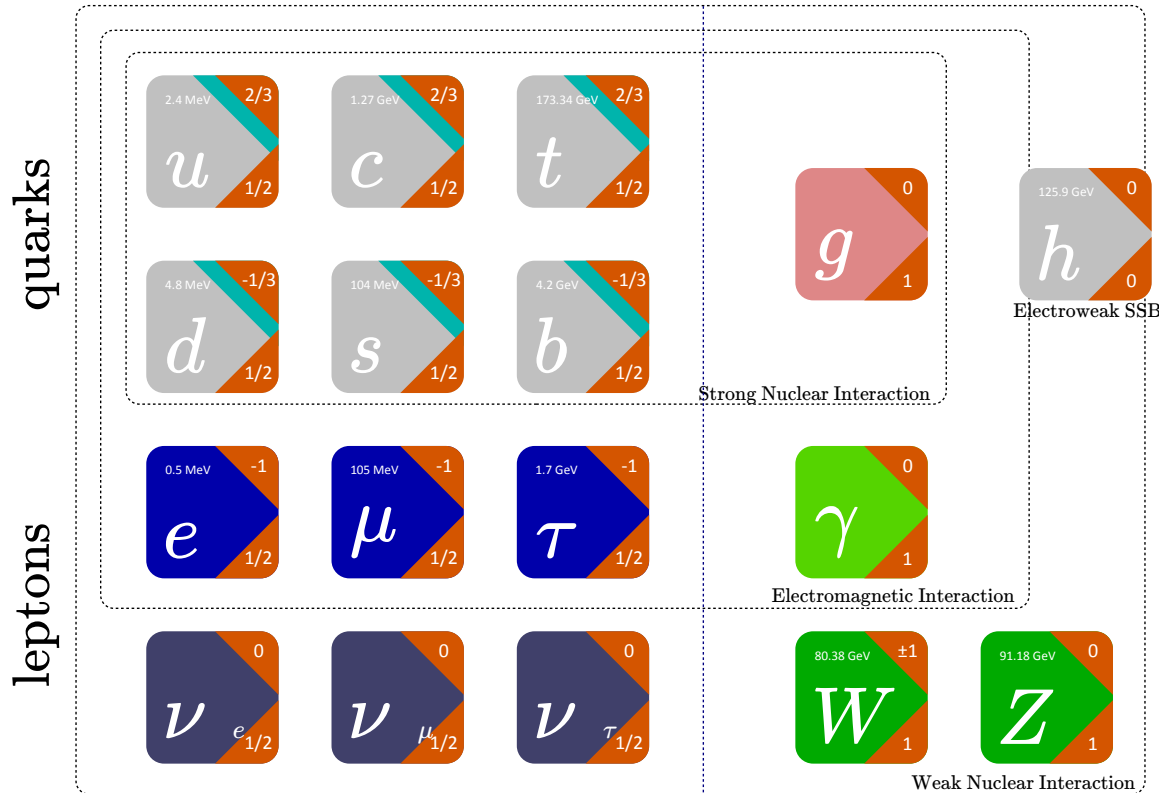
Standard Model Neutrinos only interact via weak boson: Z and W



$$\mathcal{L}_{\text{I,L}}^{(\text{NC})} = - \frac{g}{2 \cos \vartheta_{\text{W}}} \left\{ \overline{\nu_{eL}} \not{Z} \nu_{eL} - (1 - 2 \sin^2 \vartheta_{\text{W}}) \overline{e_L} \not{Z} e_L + 2 \sin^2 \vartheta_{\text{W}} \overline{e_R} \not{Z} e_R \right\} \\ + g \sin \vartheta_{\text{W}} \overline{e} \not{A} e$$

The Standard Model (fail)

SM matter families



- The Standard Model predicts massless neutrinos
- Neutrino Oscillation experiments require $O(\text{eV})$ neutrino masses

Neutrino Oscillations



Neutrino oscillations

Flavor and mass eigenstates **do not coincide** $|\nu_\alpha\rangle = \sum_k U_{\alpha k}^* |\nu_k\rangle$

Mass eigenstates
evolve according to:

$$\left| \begin{aligned} i \frac{\partial \Psi}{\partial t} &= \mathcal{H} \Psi \\ \mathcal{H}_{\text{vac}} &= \frac{1}{2E} U \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} U^\dagger \end{aligned} \right.$$

The final ν flavor depends on: **Initial state, Source distance, Neutrino energy**

With more details:

Flavor and mass eigenstates are related by a unitary transformation:

$$|\nu_\alpha\rangle = \sum_k U_{\alpha k}^* |\nu_k\rangle \quad (\alpha = e, \mu, \tau)$$

where the eigenstates satisfy:

$$\mathcal{H}|\nu_k\rangle = E_k|\nu_k\rangle \quad E_k = \sqrt{\vec{p}^2 + m_k^2}$$

and these states evolve according to the time-evolving operator:

$$i \frac{d}{dt} |\nu_k(t)\rangle = \mathcal{H} |\nu_k(t)\rangle \quad |\nu_k(t)\rangle = e^{-iE_k t} |\nu_k\rangle$$

$$A_{\nu_\alpha \rightarrow \nu_\beta}(t) \equiv \langle \nu_\beta | \nu_\alpha(t) \rangle = \sum_k U_{\alpha k}^* U_{\beta k} e^{-iE_k t}$$

With more details:

Since neutrino can be detected using weak interactions, the flavor evolution is:

$$|\nu_\alpha(t)\rangle = \sum_k U_{\alpha k}^* e^{-iE_k t} |\nu_k\rangle \qquad |\nu_\alpha(t)\rangle = \sum_{\beta=e,\mu,\tau} \left(\sum_k U_{\alpha k}^* e^{-iE_k t} U_{\beta k} \right) |\nu_\beta\rangle$$

and the transition amplitude between flavors corresponds to:

$$A_{\nu_\alpha \rightarrow \nu_\beta}(t) \equiv \langle \nu_\beta | \nu_\alpha(t) \rangle = \sum_k U_{\alpha k}^* U_{\beta k} e^{-iE_k t}$$

but the observable is the transition probability:

$$P_{\nu_\alpha \rightarrow \nu_\beta}(t) = |A_{\nu_\alpha \rightarrow \nu_\beta}(t)|^2 = \sum_{k,j} U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^* e^{-i(E_k - E_j)t}$$

$$E_k - E_j \simeq \frac{\Delta m_{kj}^2}{2E}$$

With more details:

The expression for the probability is in term of time:

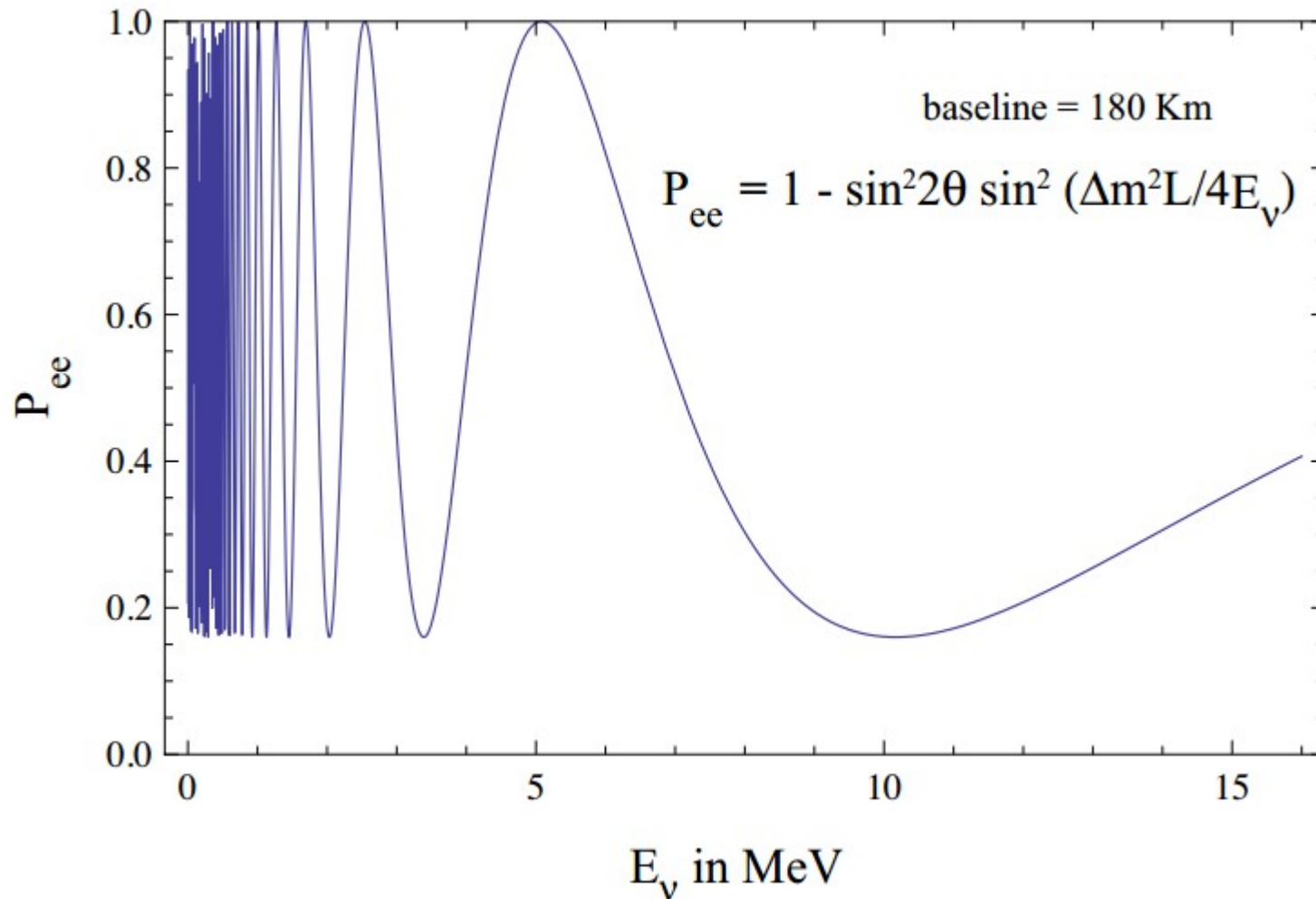
$$P_{\nu_\alpha \rightarrow \nu_\beta}(t) = \sum_{k,j} U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^* \exp \left(-i \frac{\Delta m_{kj}^2 t}{2E} \right)$$

but it is more useful in term of distance:

$$P_{\nu_\alpha \rightarrow \nu_\beta}(L, E) = \sum_{k,j} U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^* \exp \left(-i \frac{\Delta m_{kj}^2 L}{2E} \right)$$

$$U = \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{bmatrix} \\ \times \text{diag}(1, e^{i\frac{\alpha_{21}}{2}}, e^{i\frac{\alpha_{31}}{2}}) .$$

With more details:

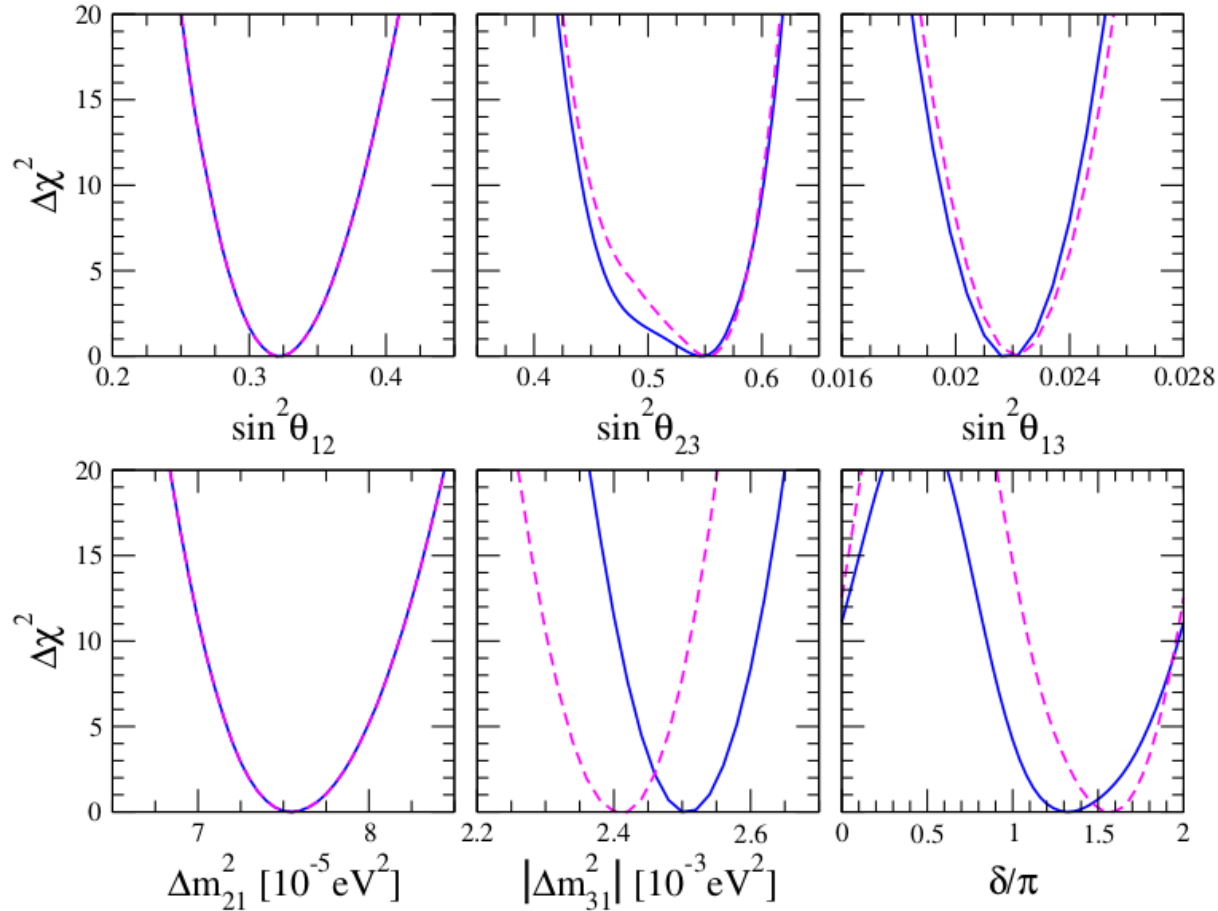


With more details:

Type of experiment	L	E	Δm^2 sensitivity
Reactor SBL	~ 10 m	~ 1 MeV	~ 0.1 eV ²
Accelerator SBL (Pion DIF)	~ 1 km	$\gtrsim 1$ GeV	$\gtrsim 1$ eV ²
Accelerator SBL (Muon DAR)	~ 10 m	~ 10 MeV	~ 1 eV ²
Accelerator SBL (Beam Dump)	~ 1 km	$\sim 10^2$ GeV	$\sim 10^2$ eV ²
Reactor LBL	~ 1 km	~ 1 MeV	$\sim 10^{-3}$ eV ²
Accelerator LBL	$\sim 10^3$ km	$\gtrsim 1$ GeV	$\gtrsim 10^{-3}$ eV ²
ATM	$20\text{--}10^4$ km	$0.5\text{--}10^2$ GeV	$\sim 10^{-4}$ eV ²
Reactor VLB	$\sim 10^2$ km	~ 1 MeV	$\sim 10^{-5}$ eV ²
Accelerator VLB	$\sim 10^4$ km	$\gtrsim 1$ GeV	$\gtrsim 10^{-4}$ eV ²
SOL	$\sim 10^{11}$ km	$0.2\text{--}15$ MeV	$\sim 10^{-12}$ eV ²

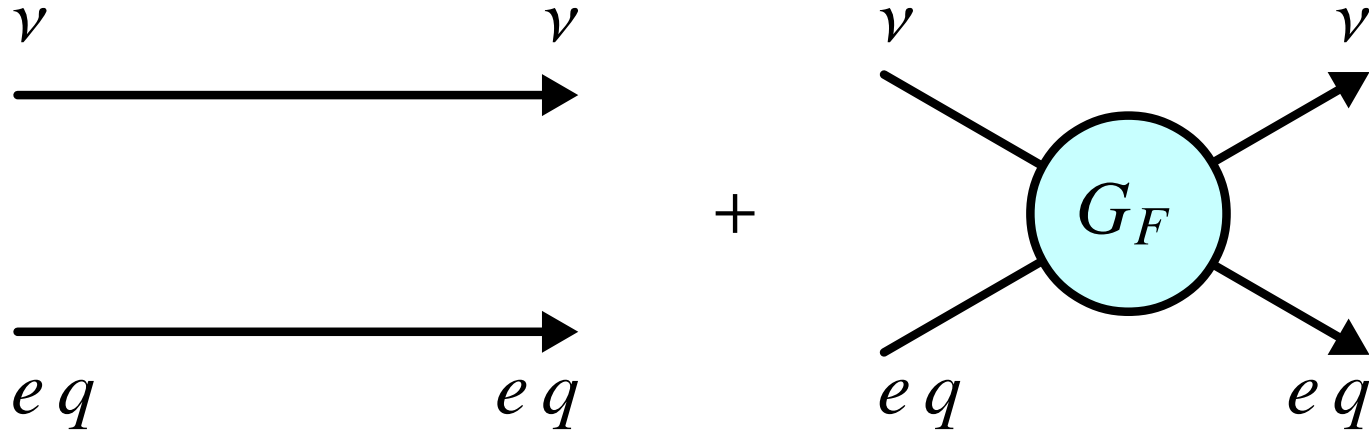
Neutrino oscillations

- The analysis of different neutrino oscillation experiments help to constraint mixing angle and mass splitting.



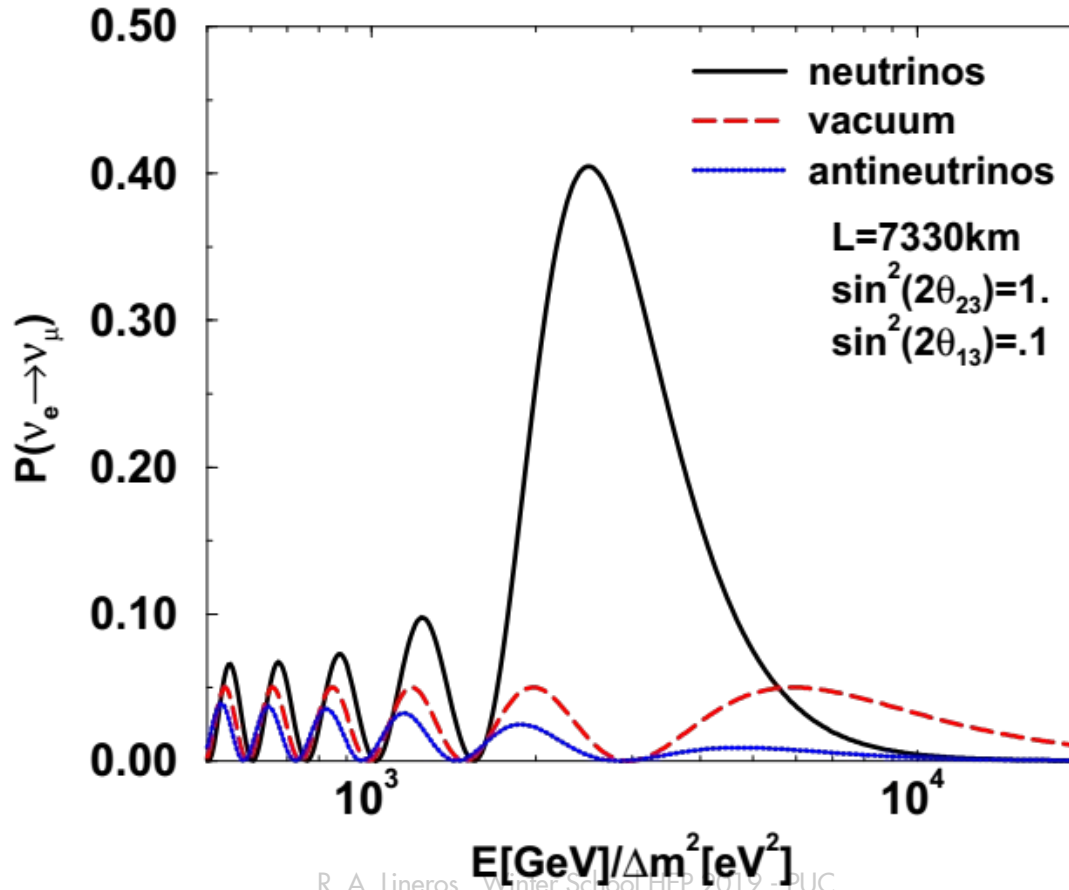
Matter effects (a.k.a. MSW effect)

The interaction with a medium modifies the oscillation patterns w.r.t. vacuum



$$\mathcal{H}_{\text{tot}} = \mathcal{H}_{\text{vac}} + \mathcal{V}$$

Matter effects (a.k.a. MSW effect)



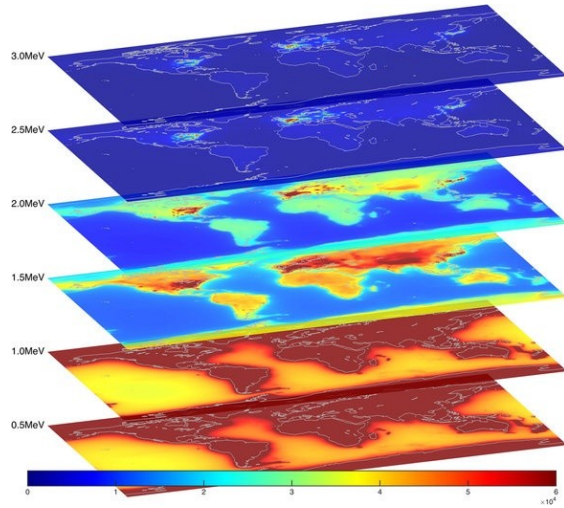
Mass Mechanisms



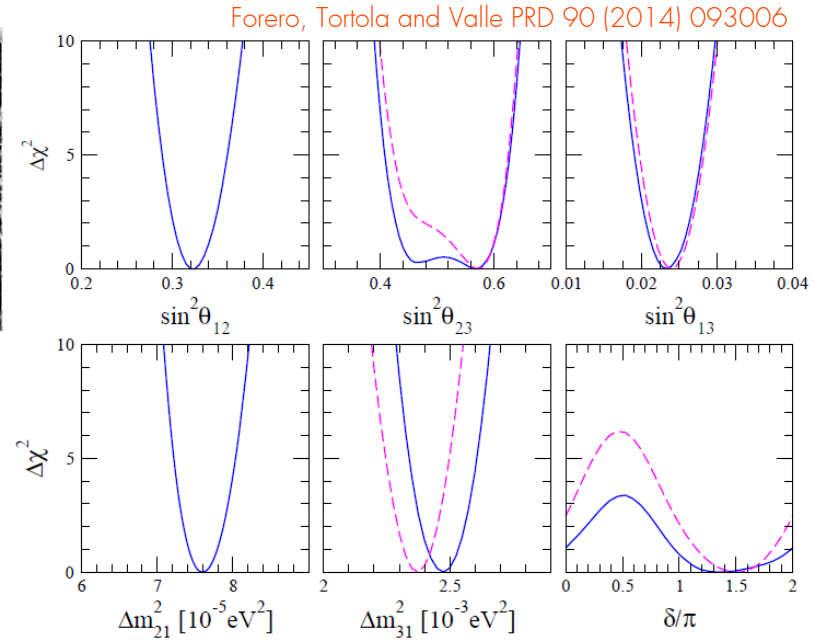
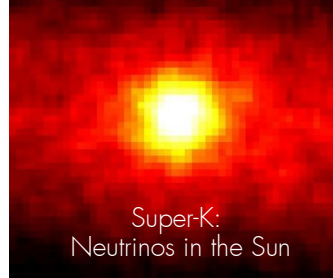
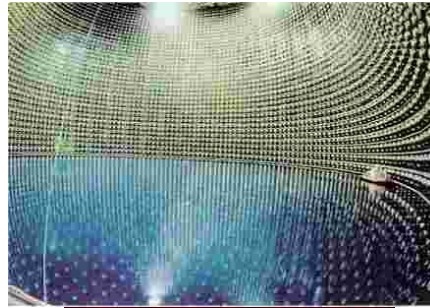
To add to the slides

Neutrino mass mechanism
Sources of neutrinos
Astroparticles neutrinos
Solar neutrinos (DM also)
Icecube neutrinos
Cosmology neutrino

Neutrinos



AGM2015: Antineutrino Global Map 2015



The **SM** predicts massless neutrinos

Physics **beyond SM** is needed to explain neutrino's mass spectrum and mixing angles

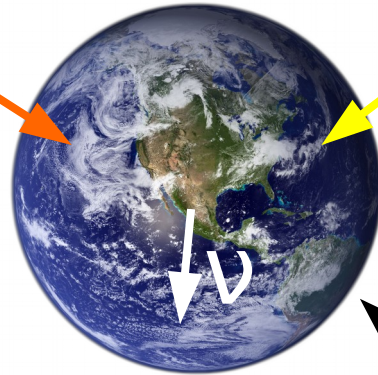
$$\mathcal{O}_{5ij} \propto (L_i H)^T (L_j H)$$

Neutrino Sources



Sun

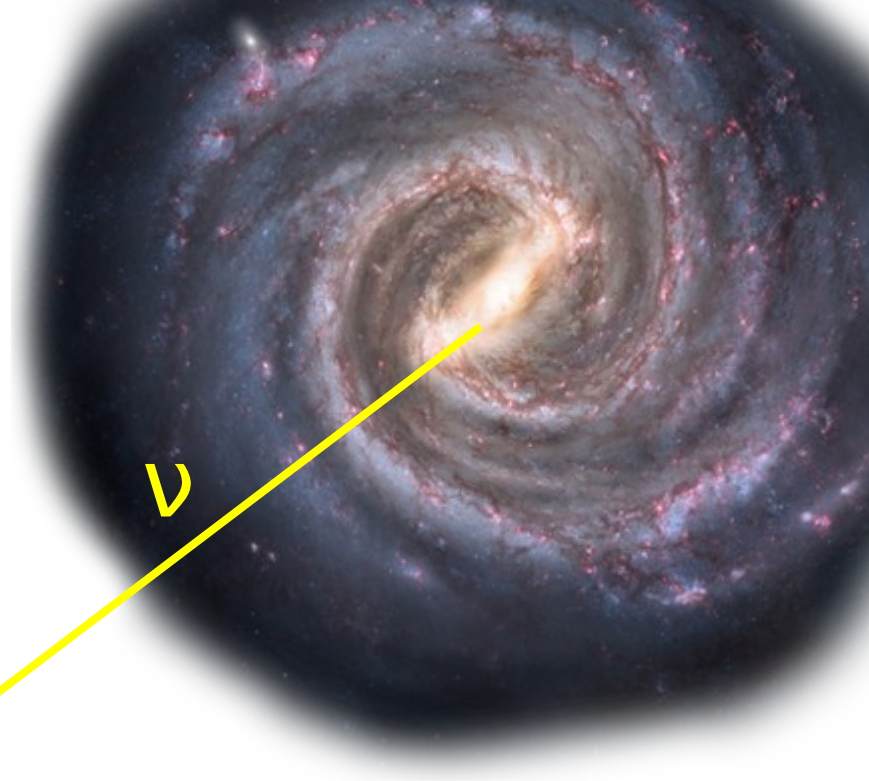
ν



Geoneutrinos
Atmospheric

ν

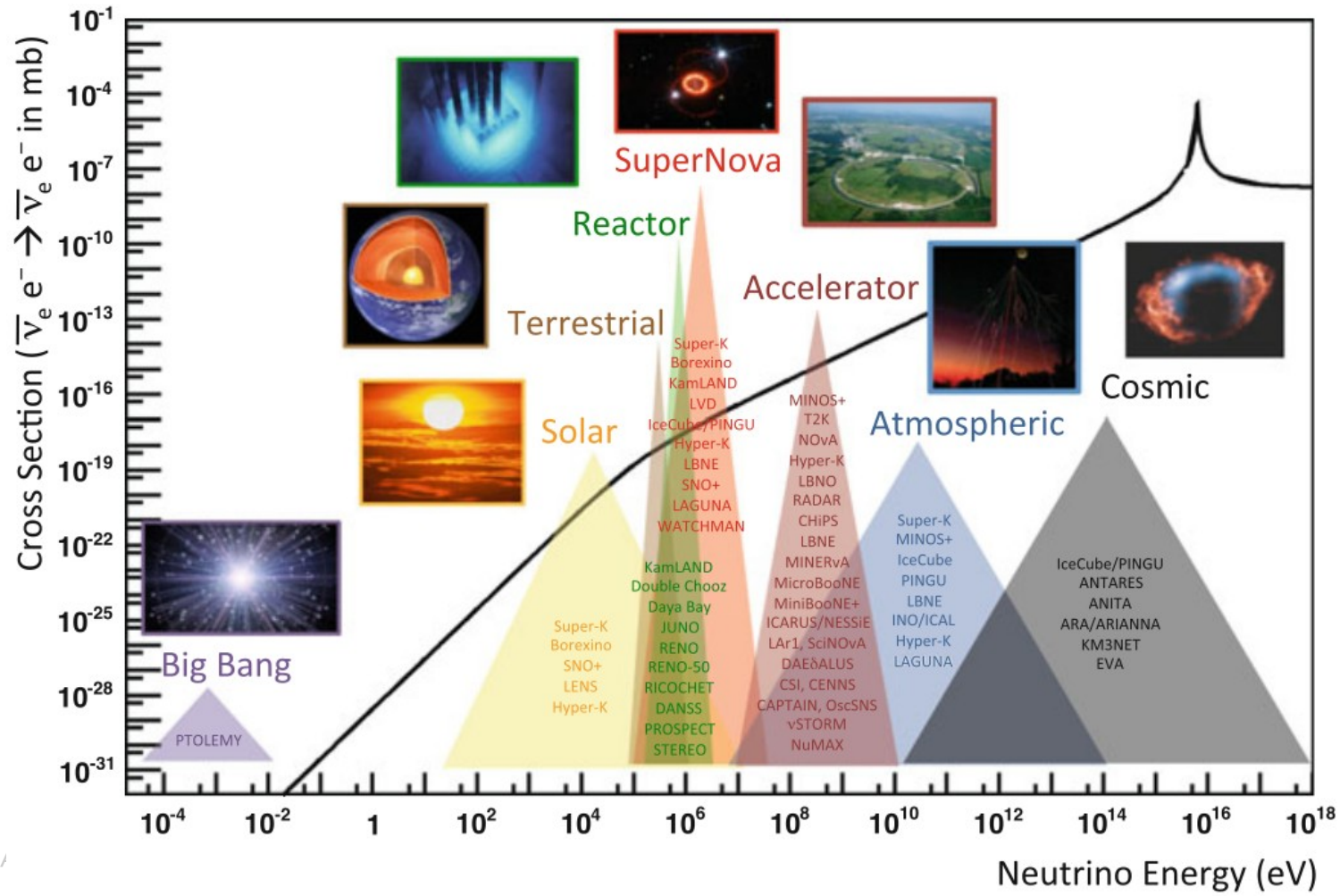
ν



ν

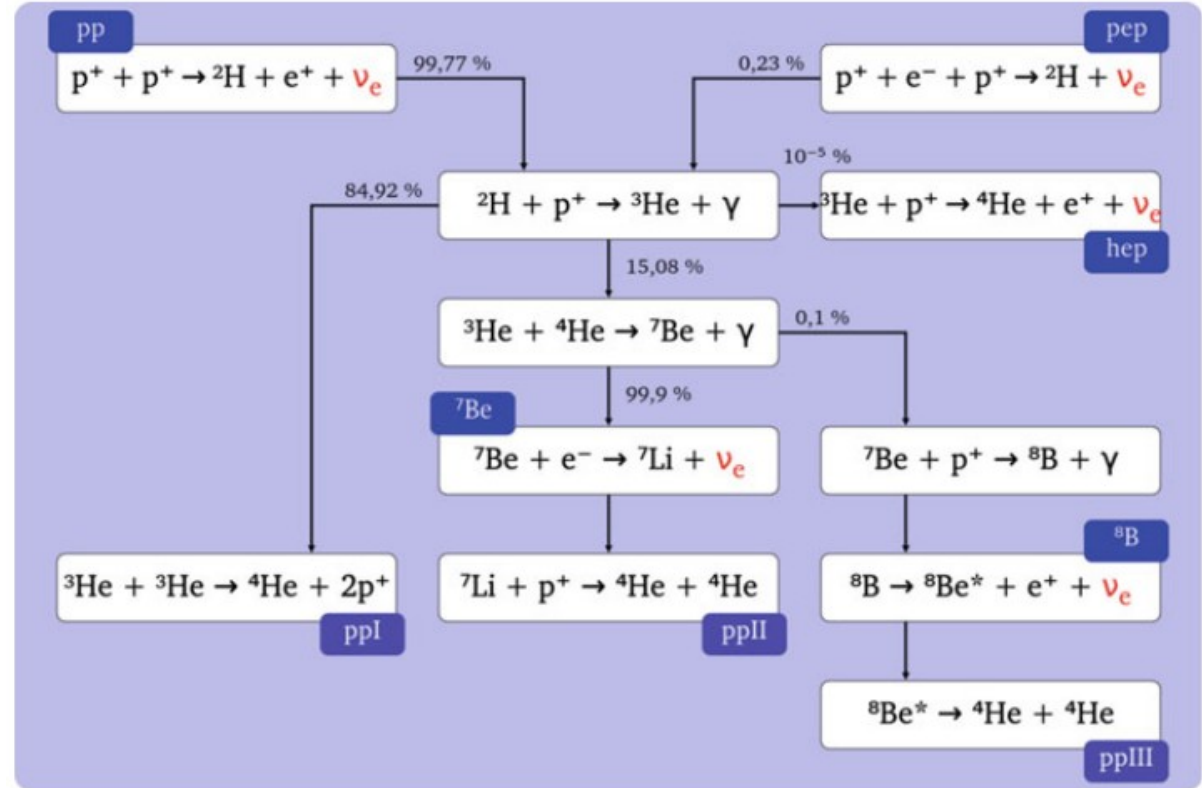
Milky Way halo
Galactic Center
Satellite galaxies
AGN

Neutrino Sources



Neutrinos from the Sun

The Sun is the brightest source of neutrinos at the MeV scale



Neutrinos from the Sun

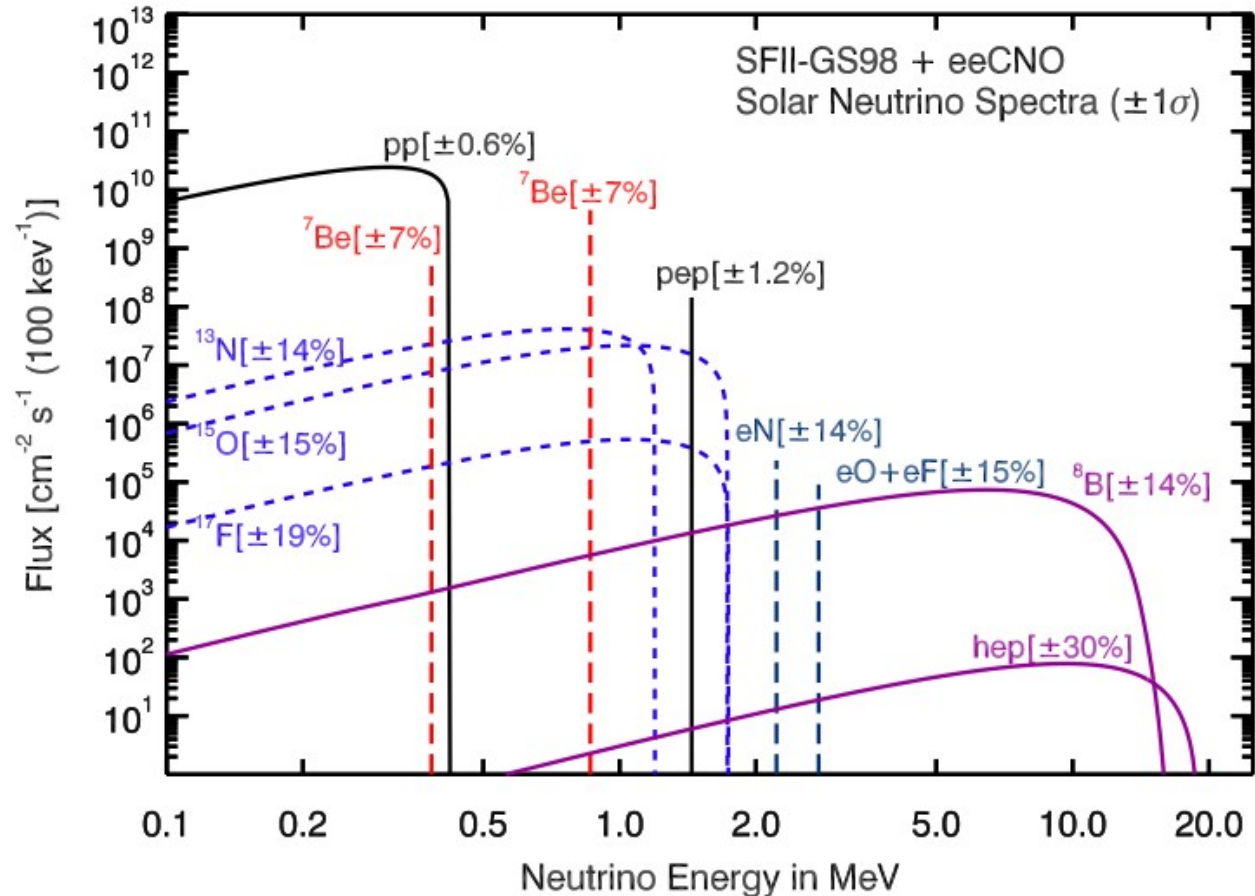
The Sun is the brightest source of neutrinos at the MeV scale



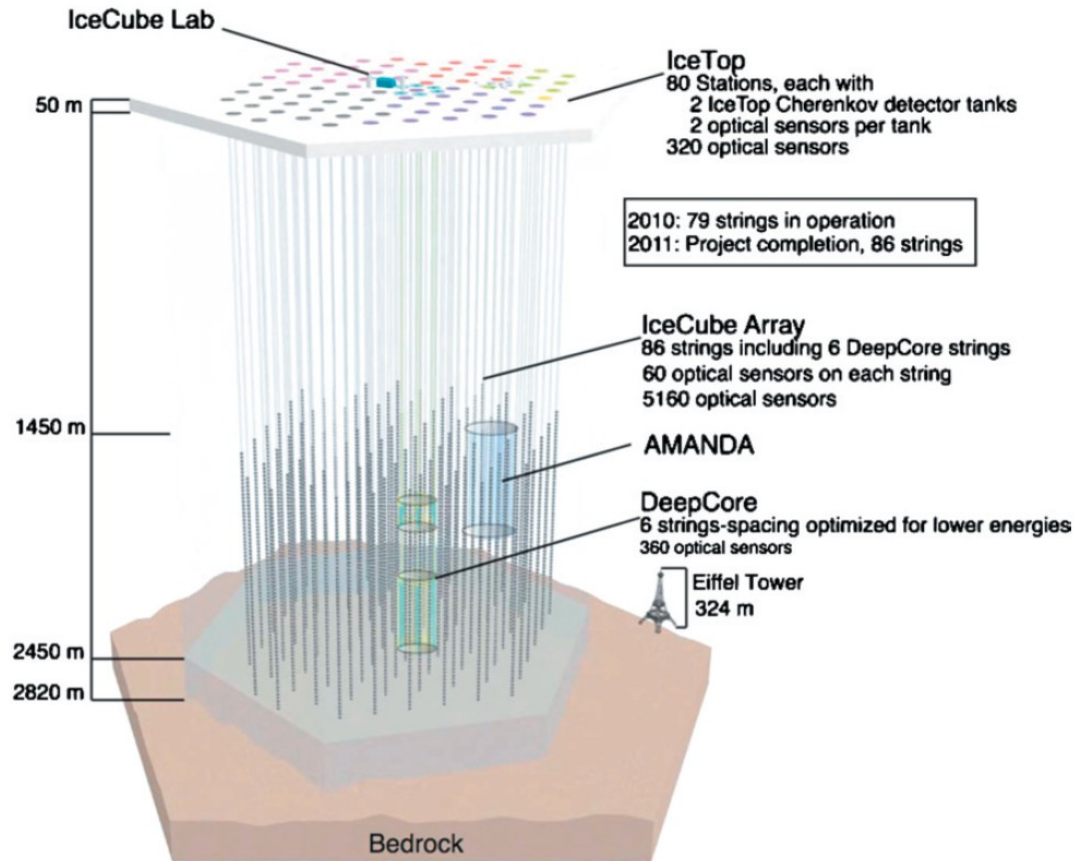
Reaction	Abbr.	Flux ($\text{cm}^{-2} \text{s}^{-1}$)
$pp \rightarrow d e^+ \nu$	pp	$5.98(1 \pm 0.006) \times 10^{10}$
$pe^- p \rightarrow d \nu$	pep	$1.44(1 \pm 0.01) \times 10^8$
${}^3\text{He } p \rightarrow {}^4\text{He } e^+ \nu$	hep	$7.98(1 \pm 0.30) \times 10^3$
${}^7\text{Be } e^- \rightarrow {}^7\text{Li } \nu + (\gamma)$	${}^7\text{Be}$	$4.93(1 \pm 0.06) \times 10^9$
${}^8\text{B} \rightarrow {}^8\text{Be}^* e^+ \nu$	${}^8\text{B}$	$5.46(1 \pm 0.12) \times 10^6$
${}^{13}\text{N} \rightarrow {}^{13}\text{C } e^+ \nu$	${}^{13}\text{N}$	$2.78(1 \pm 0.15) \times 10^8$
${}^{15}\text{O} \rightarrow {}^{15}\text{N } e^+ \nu$	${}^{15}\text{O}$	$2.05(1 \pm 0.17) \times 10^8$
${}^{17}\text{F} \rightarrow {}^{17}\text{O } e^+ \nu$	${}^{17}\text{F}$	$5.29(1 \pm 0.20) \times 10^6$

Neutrinos from the Sun

The Sun is the brightest source of neutrinos at the MeV scale



Neutrino Observatories



IceCube (located in Antarctica)



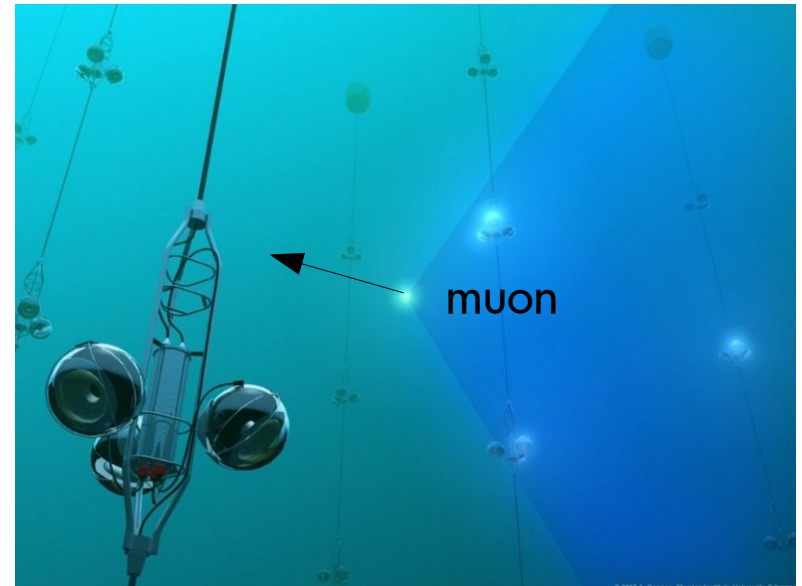


Neutrino Observatories



ANTARES

(located in the mediterranean sea)

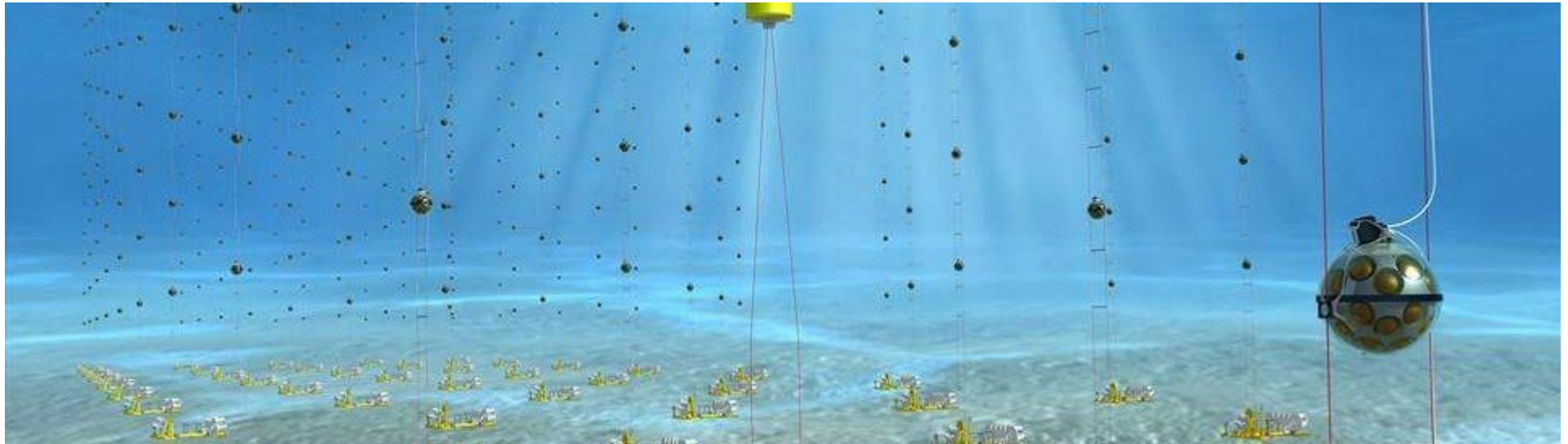


Neutrino Observatories



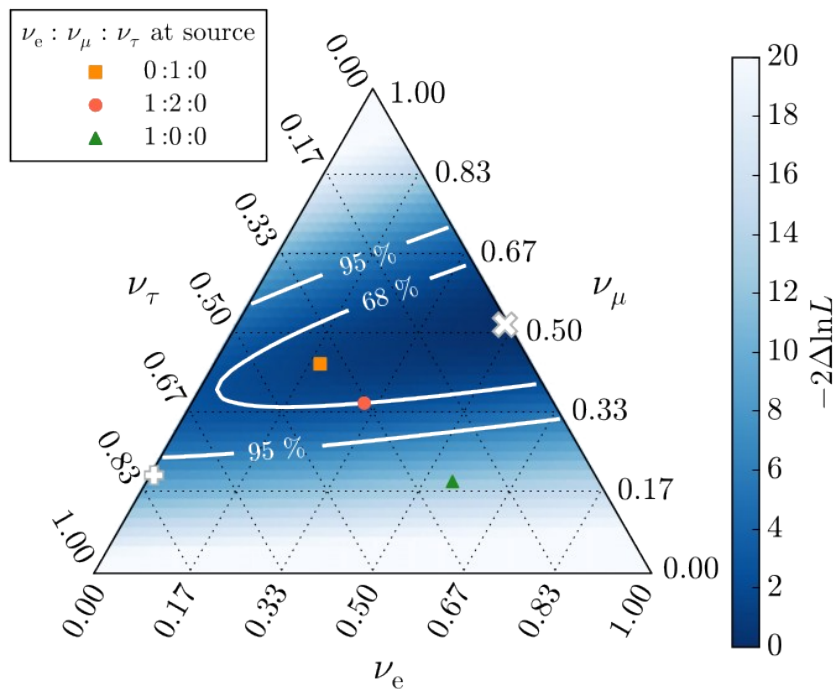
KM3NET

(located in the mediterranean sea)



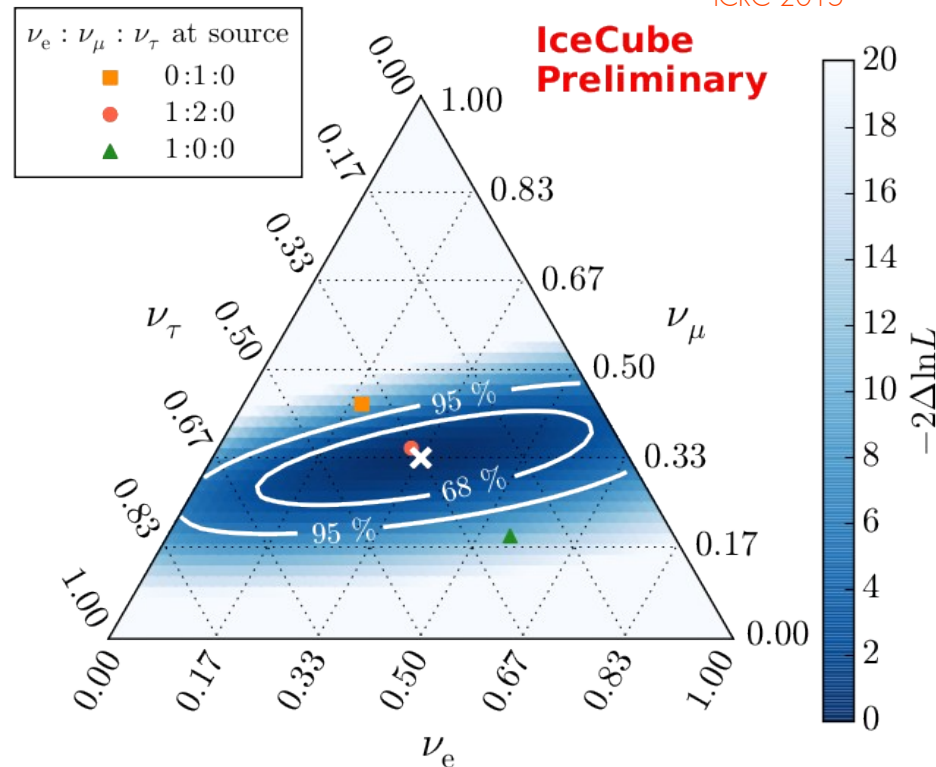
Flavor composition in IceCube

THE ASTROPHYSICAL JOURNAL, 809:98 (15pp), 2015 August 10



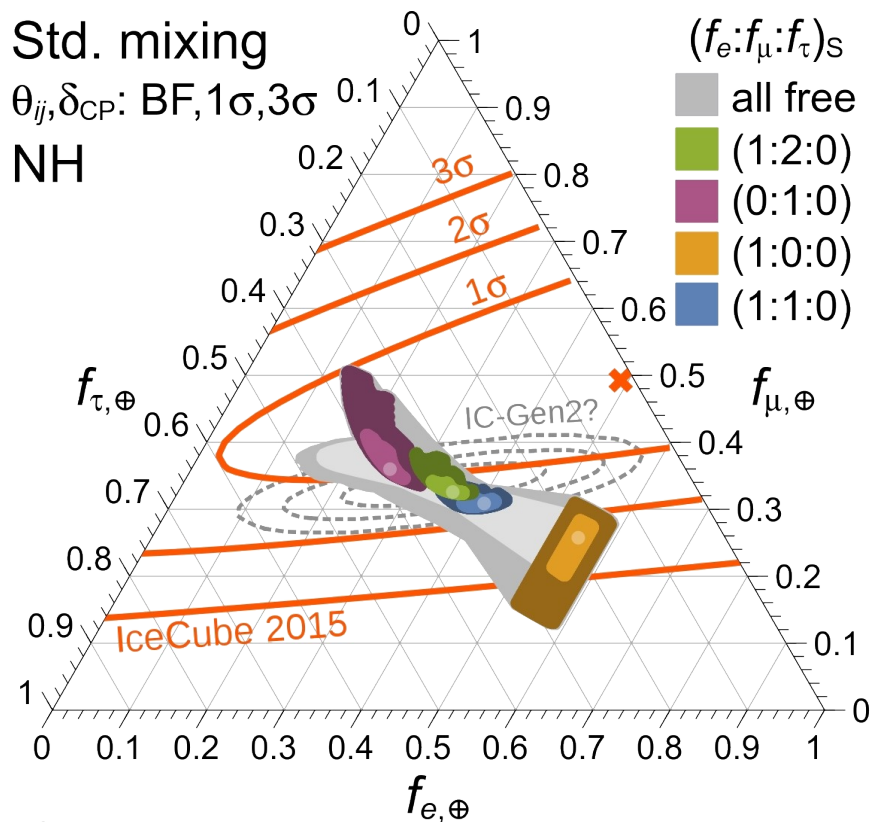
Latest result on flavor composition

ICRC 2015

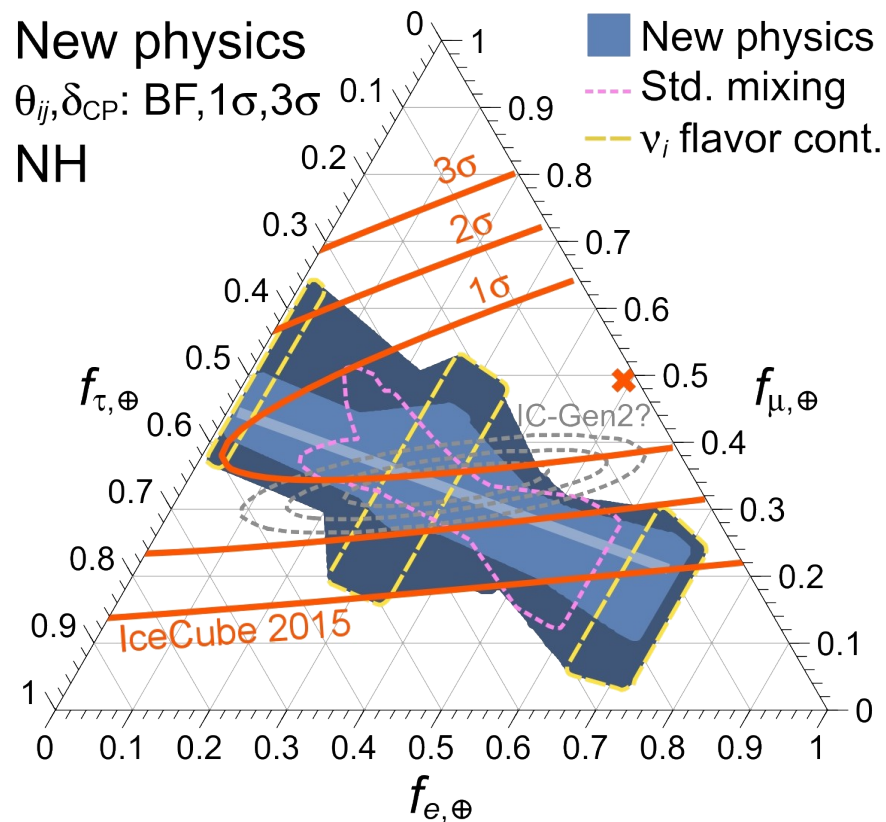


Sensitivity after 10 years

Flavor composition

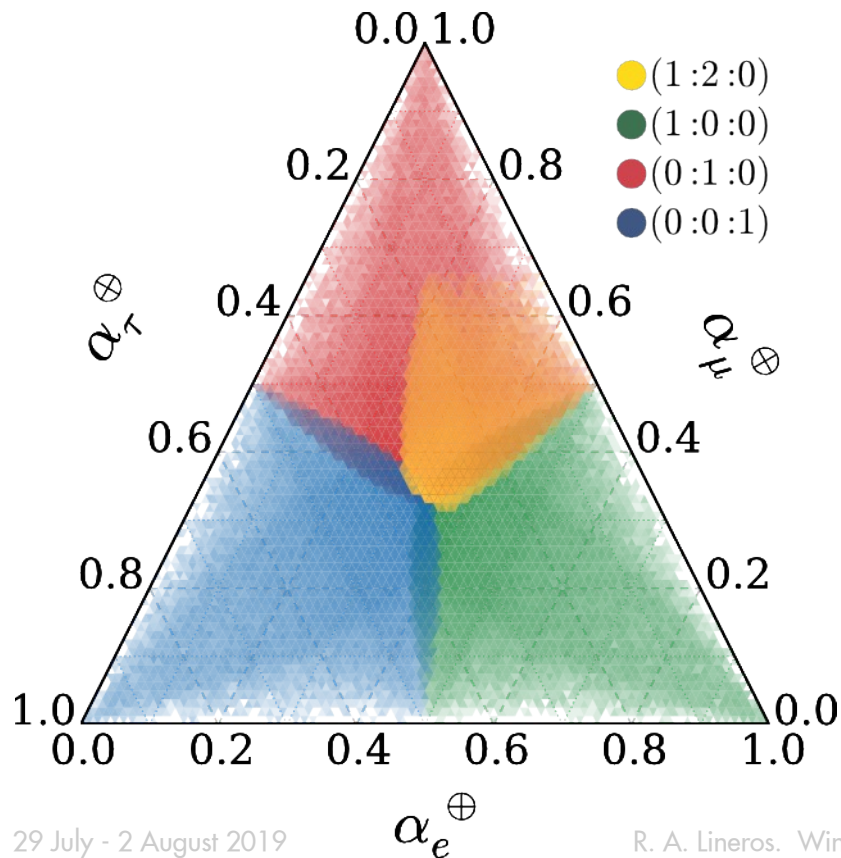


M. Bustamante et al. PRL 115, 161302 (2015)



Effects from New Physics

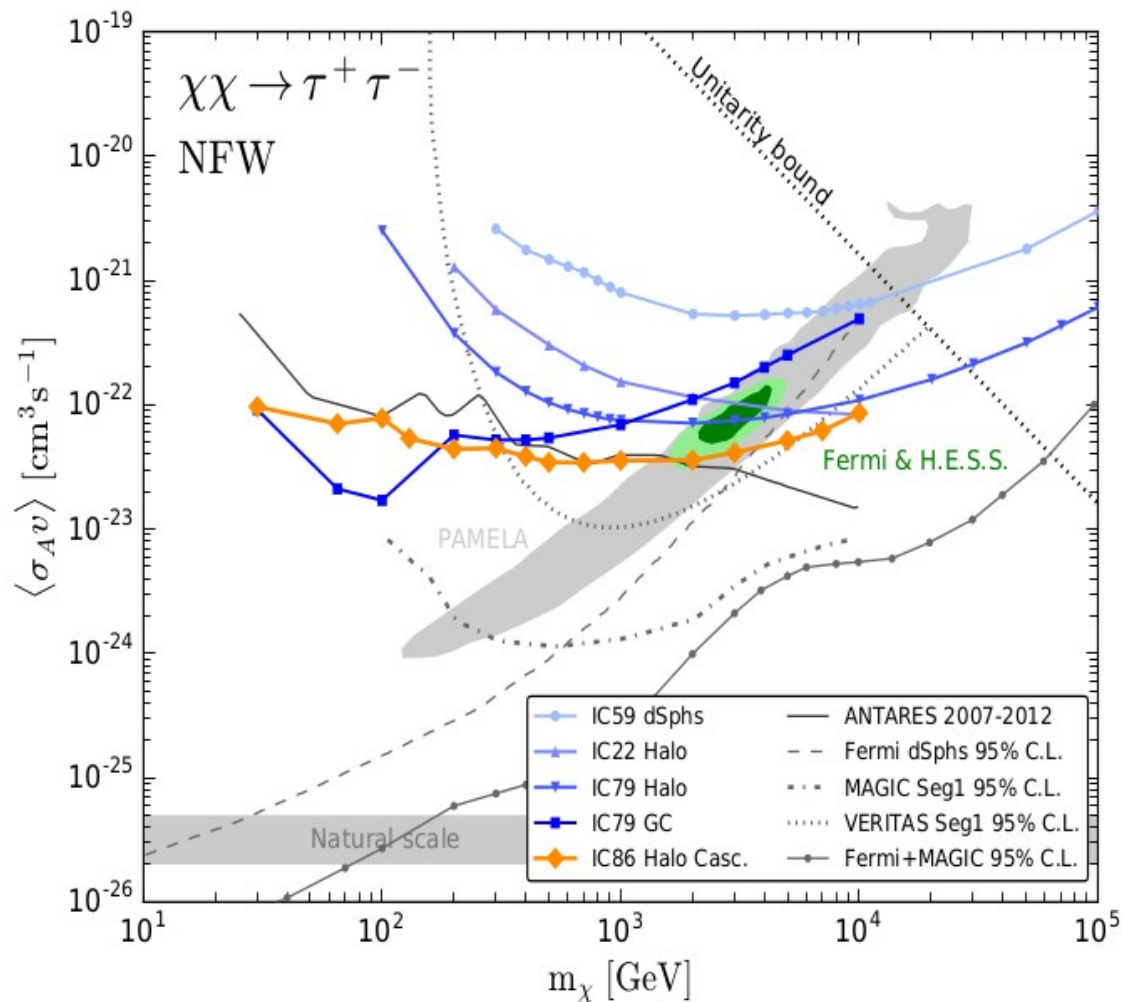
Argüelles et al. PRL 115, 161303 (2015)



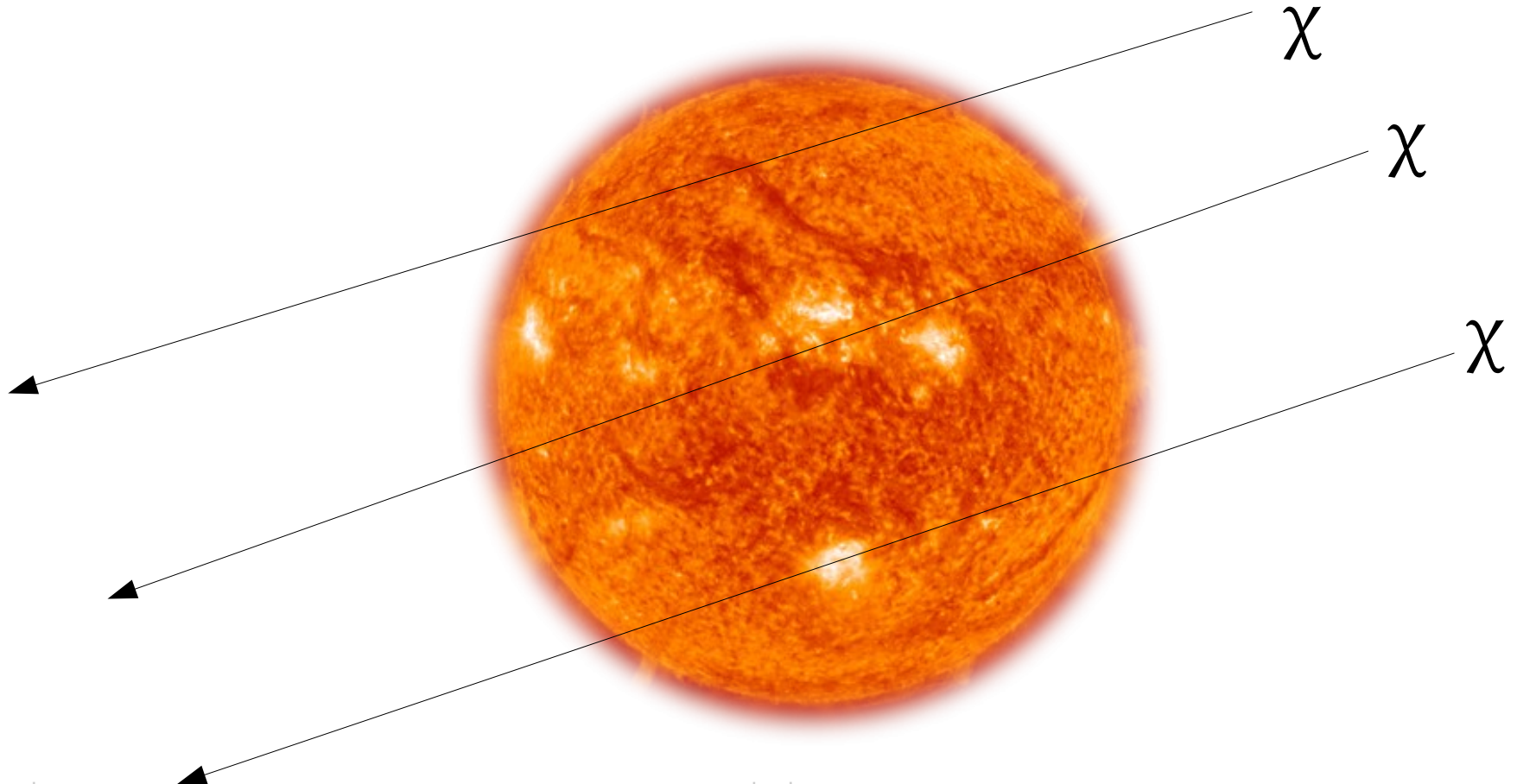
Sources of New Physics:

- Space torsion
- CPT - Lorentz violation
- Dark Matter interactions

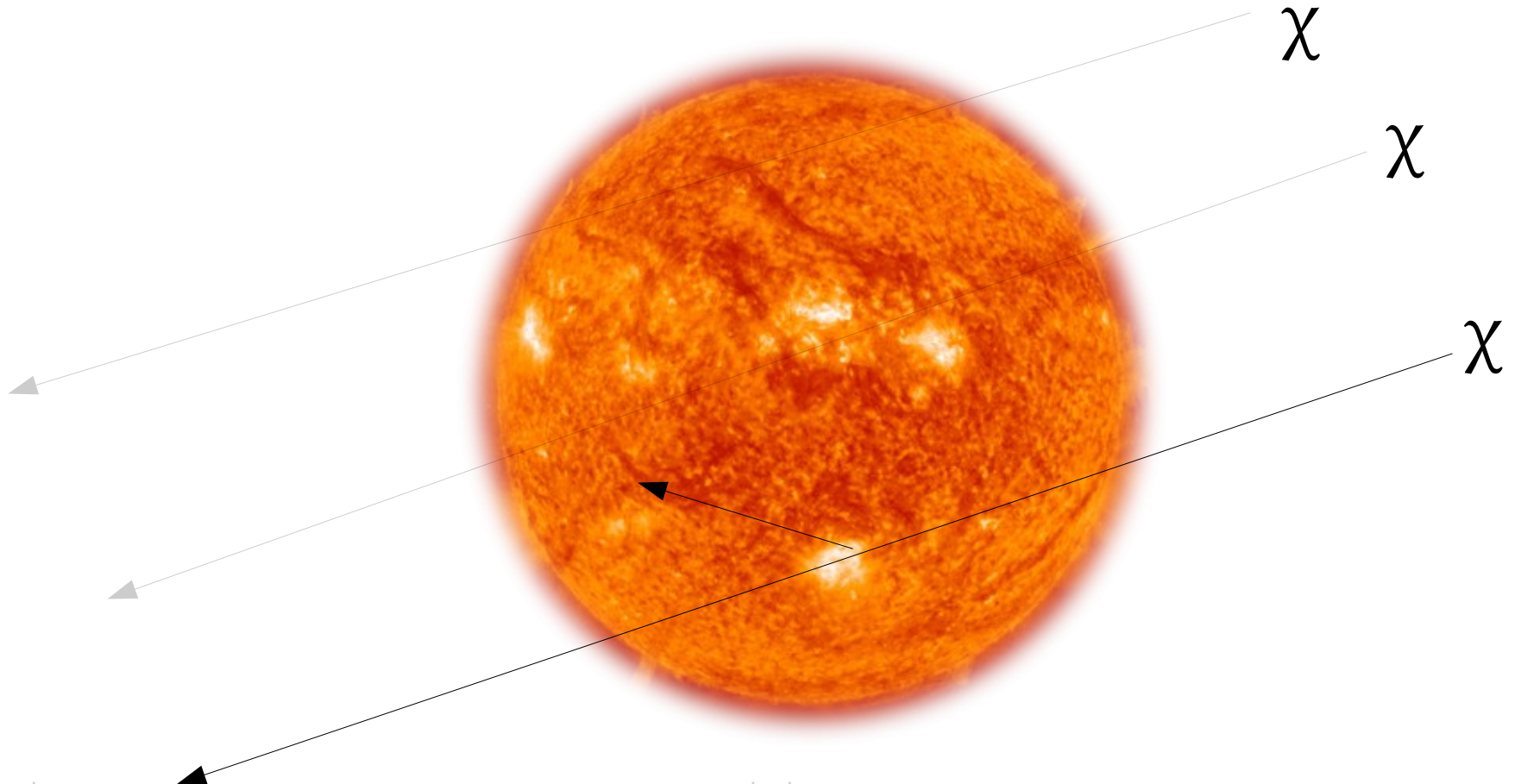
Dark Matter Searches



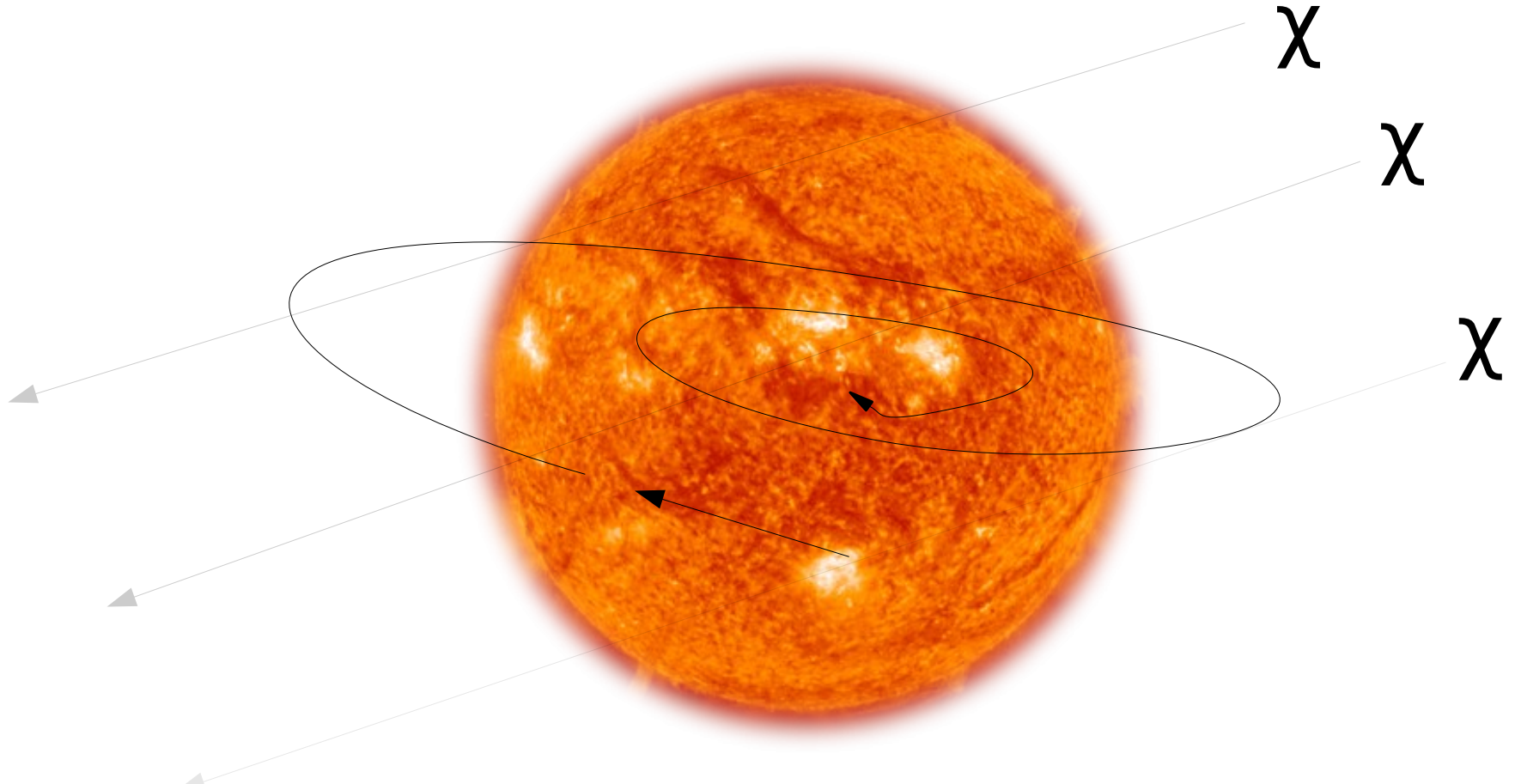
Dark Matter in the Sun



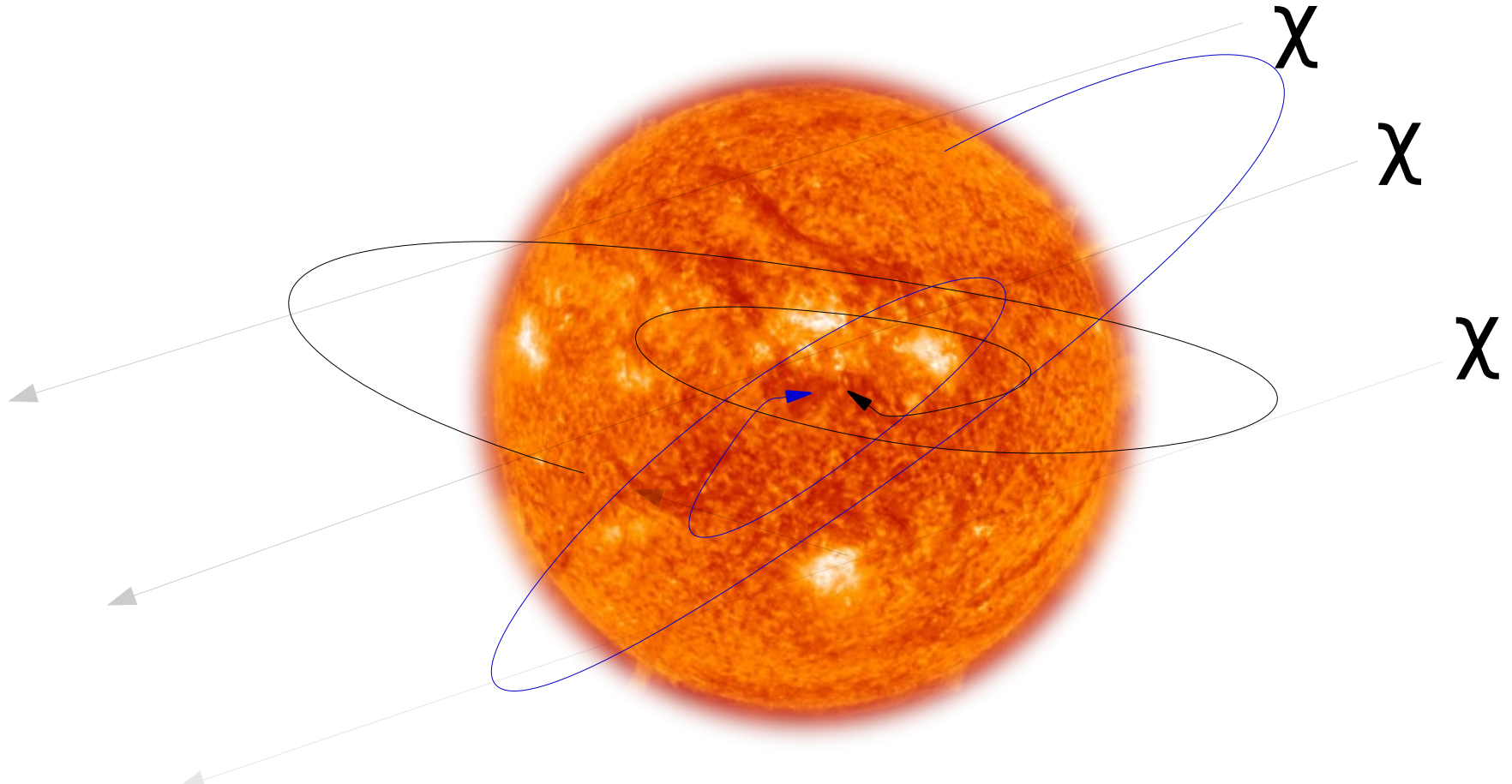
Dark Matter in the Sun



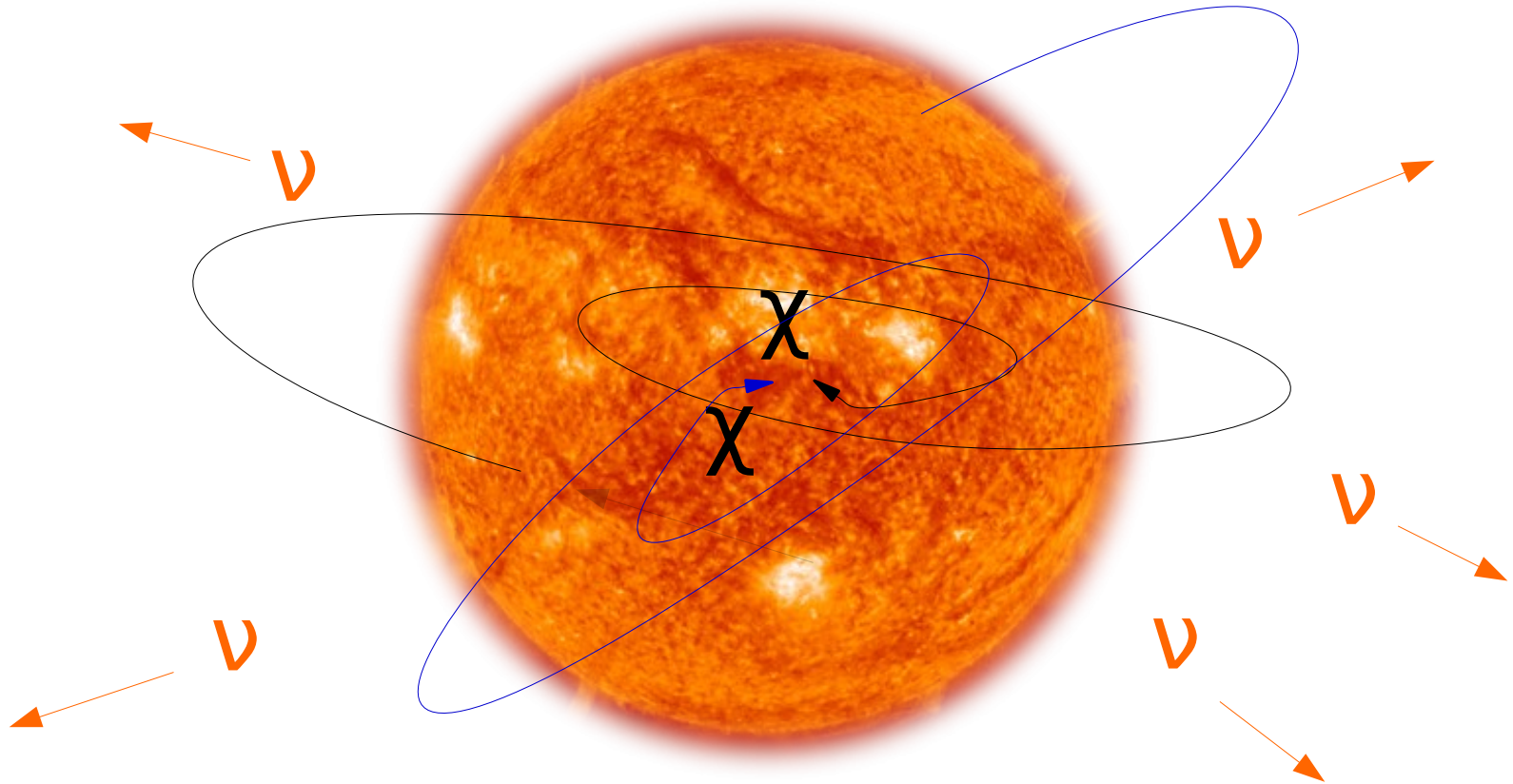
Dark Matter in the Sun



Dark Matter in the Sun

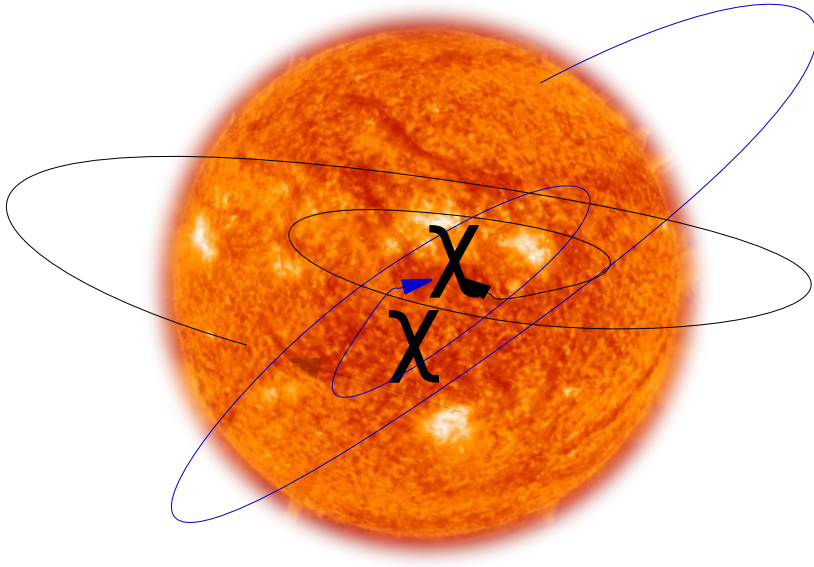


Dark Matter in the Sun



Dark Matter in the Sun

arxiv:1005.5711



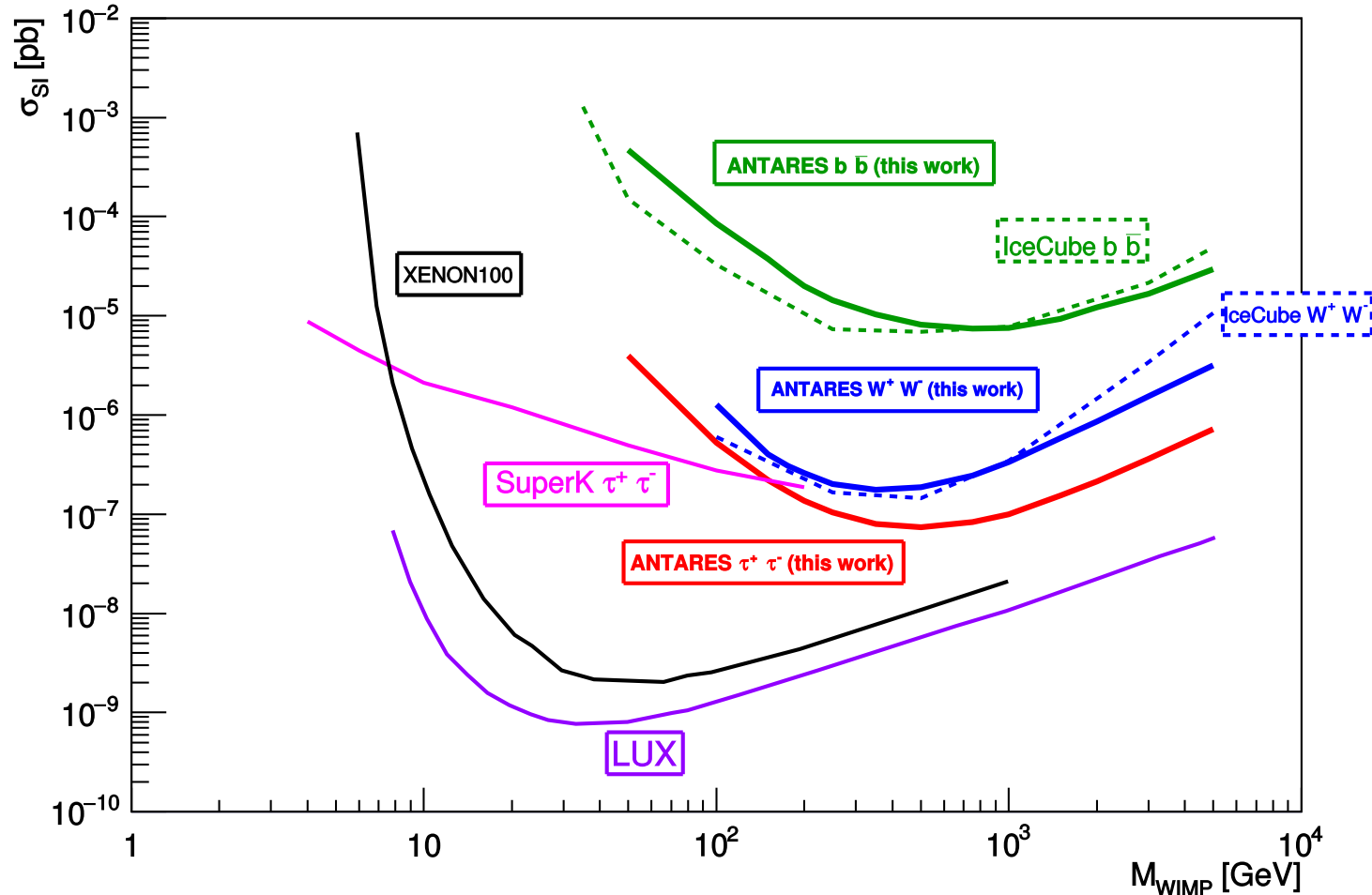
$$\dot{N}_\chi = C - 2AN_\chi^2 - EN_\chi$$

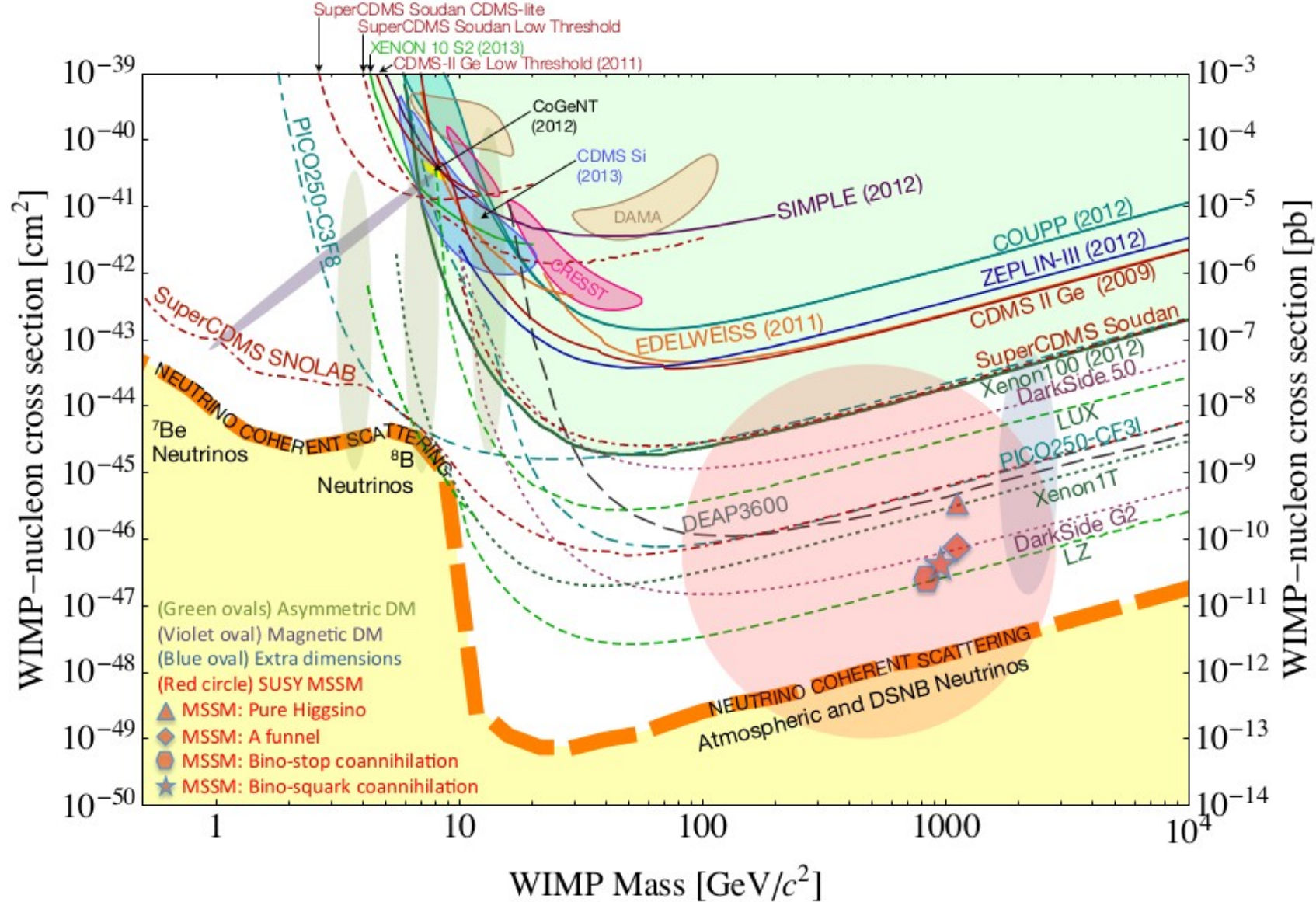
Capture

Annihilation

Evaporation

Dark Matter in the Sun





The Tools



Tomorrow we continue...

Sterile
neutrino

Electron
neutrino

Muon
neutrino

Tau
neutrino

