

Supplementary Information

Self Assembly of Ultrathin Au Nanowires into Expanded Hexagonal Superlattice Studied by in situ SAXS

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1. Model for the SAXS intensity p2
2. Supplementary images p6
 - Figure S1 : TEM images of nanoparticles obtained at 40°C under magnetic stirring after 5 et 24h
 - Figure S2 : TEM images of nanoparticles obtained after 5h at 40°C for different OY concentrations
 - Figure S3: TEM image of Au NPs obtained after 27h at 40°C for TIPS/Au = 100 and [OY] = 400 mM.
 - Figure S4 : XPS survey spectrum of purified Au NWs
 - Figure S5 : High resolution XPS spectra at Au and N edges
 - Figure S6 : ^1H MAS spectra of purified Au NWs
 - Figure S7 : TEM images of kinetic study at 25°C
 - Figure S8 : SAXS spectra of the kinetic study at 25 °C
 - Figure S9 : Volume fraction evolution as a function of time at 25 °C
 - Figure S10: TEM image of assembled Au NWs.

1. Model for the SAXS intensity.

The SAXS results after background subtraction have been interpreted using a model with three contributions for the intensity (measured in absolute scale, cm^{-1} units): spherical gold nanoparticles, gold nanowires and a hexagonal phase of nanowires, plus a small remaining background constant C :

$$I(q) = I_{\text{sphere}}(q) + I_{\text{NW}}(q) + I_{\text{Bragg}}(q) + C$$

All data have been modeled in a similar way as in references [1-5], in the context of the self-assembly of silica nano-structured materials, using a least-square fitting procedure. Typical values of C less than 0.02 cm^{-1} were used in the fits.

- Spheres contribution:

A polydisperse form factor of spheres was used, with no structure factor contribution because of the very low overall volume fraction:

$$I_{\text{sphere}}(q) = n_s (\rho_{\text{gold}} - \rho_0)^2 P_{\text{sphere}}(q)$$

where n_s is the number density of the spheres and ρ_{gold} and ρ_0 the scattering length densities of gold and hexane. For SAXS, the scattering length density is the product of the classical radius of the electron $r_e = 2.81794 \times 10^{-15} \text{ m}$ by the electron density (number of electrons per unit volume). The values are known from literature: $\rho_{\text{gold}} = r_e \rho_{\text{gold}}^e$ and $\rho_0 = r_e \rho_0^e$, with $\rho_{\text{gold}}^e = 4650 \text{ e nm}^{-3}$ and $\rho_{\text{hexane}}^e = 230 \text{ e nm}^{-3}$.

The amplitude form factor of a sphere of radius R is $f_s(q, R)$:

$$f_s(q, R) = 3(\sin(qR) - qR \cos(qR))/(qR)^3$$

and $V_s = (4\pi/3)R^3$.

The averaged spherical form factor is:

$$P_{sphere}(q) = \langle F_s^2(q) \rangle = \langle V_s^2 f_s^2(q) \rangle$$

where $\langle \rangle$ stands for the average over the polydispersity (parameter σ) using a Schulz-Zimm distribution:

$$\langle F_s^2(q, R) \rangle = \int p(r, R, \sigma) F_s^2(q, r) dr$$

$$p(r, R, \sigma) = f_{SZ}(r, R, 1/\sigma^2 - 1)$$

$$f_{SZ}(x, x_0, z) = \frac{x^z}{\Gamma(z+1)} \left(\frac{z+1}{x_0} \right)^{z+1} \exp\left(- (z+1) \frac{x}{x_0}\right)$$

The volume fraction ϕ_s of spheres in solution was derived as follows: $n_s = \phi_s / \langle V_s \rangle$

$$I_{sphere}(q) = (\rho_{gold} - \rho_0)^2 \phi_s \frac{\langle V_s^2 f_s^2(q, R) \rangle}{\langle V_s \rangle}$$

- Nanowires contribution:

A form factor for wires of constant length L and average radius R_{NW} was used, with no structure factor contribution:

$$I_{NW}(q) = n_{NW} (\rho_{gold} - \rho_0)^2 P_{NW}(q)$$

where n_{NW} is the number density of the nanowires and ρ_{gold} and ρ_0 the scattering length densities of gold and hexane, like for the spheres.

Polydispersity was introduced only for the radius R_{NW} and the length was fixed to $L = 100$ nm. Because of the large aspect ratio of the nanowires, the form factor could be factorized in two terms, where $\langle \rangle$ stands for the average over the polydispersity (parameter σ) using a Schulz-Zimm distribution as for the spheres.

$$P_{NW}(q) = P_{rod}(q) < F_{CS}^2(q) >$$

where $F_{CS}(q)$ is the form factor of the NW's cross-section of radius R_{NW}

$$F_{CS}(q) = \pi R_{NW}^2 \frac{2J_1(qR_{NW})}{qR_{NW}}$$

$$P_{rod}(q) = L^2 \left(\frac{2 \sin I(qL)}{qL} - \frac{4 \sin^2(\frac{qL}{2})}{(qL)^2} \right) \text{ with } \sin I(x) = \int_0^x \frac{\sin u}{u} du$$

For $q \gg 1/L$, $P_{rod}(q)$ exhibits the characteristic q^{-1} power low dependence:

$$P_{rod}(q) \cong \frac{\pi L}{q}$$

The volume fraction ϕ_{NW} of nanowires in solution was derived as follows:

$$n_{NW} = \phi_{NW} / < V_{NW} >$$

$$I_{NW}(q) = (\rho_{gold} - \rho_0)^2 \phi_{NW} \frac{P_{rod}(q) < F_{CS}^2(q) >}{L \pi < R_{NW}^2 >}$$

$$I_{NW}(q) = (\rho_{gold} - \rho_0)^2 \phi_{NW} \frac{< F_{CS}^2(q) >}{q < R_{NW}^2 >}$$

- Hexagonal phase:

The contribution was modeled by the powder average of a long-range 2D-hexagonal lattice of nanowires:

$$I_{Bragg}(q) = n (\rho_{gold} - \rho_0)^2 P_{NW}(q) [1 - \beta(q)G(q) + \beta(q)G(q)Z_0(q)]$$

$$Z_0(q) = \frac{1}{s_{cell}} \frac{2\pi}{q} \sum_{hk} m_{hk} L_{hk}(q, q_{hk})$$

where s_{cell} is the area of the 2D-hexagonal unit cell and m_{hk} is the multiplicity of each Bragg peak.

The same form factor $P_{NW}(q)$ was used as for the nanowires in solution. The Bragg peaks were recorded at discrete positions q_{hk} . The Bragg peaks were modeled using peak-shaped functions $L_{hk}(q, q_{hk})$ with Gaussian/Lorentzian profiles with a width inversely proportional to the average domain size D of the hexagonal domains, corresponding to the diameter of the fiber/bundle around 70 nm. The essential contribution to the Bragg intensities was the form factor $P_{NW}(q)$ of the gold nanowires. But in addition, two types of disorder were included in the model. First a small polydispersity over the nanowires radius (parameter $\sigma = 0.012$), through the average over the cross-section form factor, term $\beta(q)$, . Second, a positional disorder was introduced through the Debye-Waller term $G(q)$:

$$\beta(q) = \frac{\langle F_{CS}(q) \rangle^2}{\langle F_{CS}^2(q) \rangle}$$

$$G(q) = \exp(-q^2 \sigma_{DW}^2 a^2)$$

Typical values of $\sigma_{DW} = 0.08$ were used for the relative mean-square displacement in the Debye-Waller term.

References:

- [1] Sundblom, A.; Oliveira, C. L. P.; Palmqvist, A. E. C.; Pedersen, J. S. *J. Phys. Chem. C* **2009**, *113*, 7706.
- [2] Manet, S.; Lecchi, A.; Imperor-Clerc, M.; Zholobenko, V.; Durand, D.; Oliveira, C. L. P.; Pedersen, J. S.; Grillo, I.; Meneau, F.; Rochas, C. *J Phys Chem B* **2011**, *115*, 11318.
- [3] Manet, S.; Schmitt, J.; Imperor-Clerc, M.; Zholobenko, V.; Durand, D.; Oliveira, C. L. P.; Pedersen, J. S.; Gervais, C.; Baccile, N.; Babonneau, F.; Grillo, I.; Meneau, F.; Rochas, C. *J Phys Chem B* **2011**, *115*, 11330
- [4] Michaux, F.; Baccile, N.; Imp  rator-Clerc, M.; Malfatti, L.; Folliet, N.; Gervais, C.; Manet, S.; Meneau, F.; Pedersen, J. S.; Babonneau, F. *Langmuir* **2012**, *28*, 17477-17493 .
- [5] Schmitt, J.; Imp  rator-Clerc, M. ; Michaux, F. ; Blin, J. L.; St  b  , M. J.; Pedersen, J.S. ; Meneau, F. *Langmuir* **2013**, *29*, 2007–2023.

3. Supplementary figures

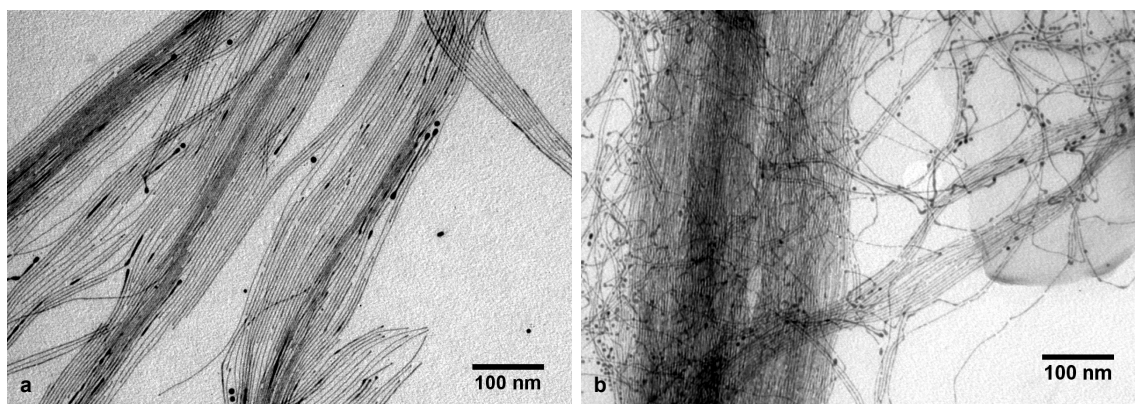


Figure S1: TEM image of Au NPs obtained after 3h at 40°C for TIPS/Au = 100 and [OY] = 400 for a) undisturbed solution or b) under magnetic stirring.

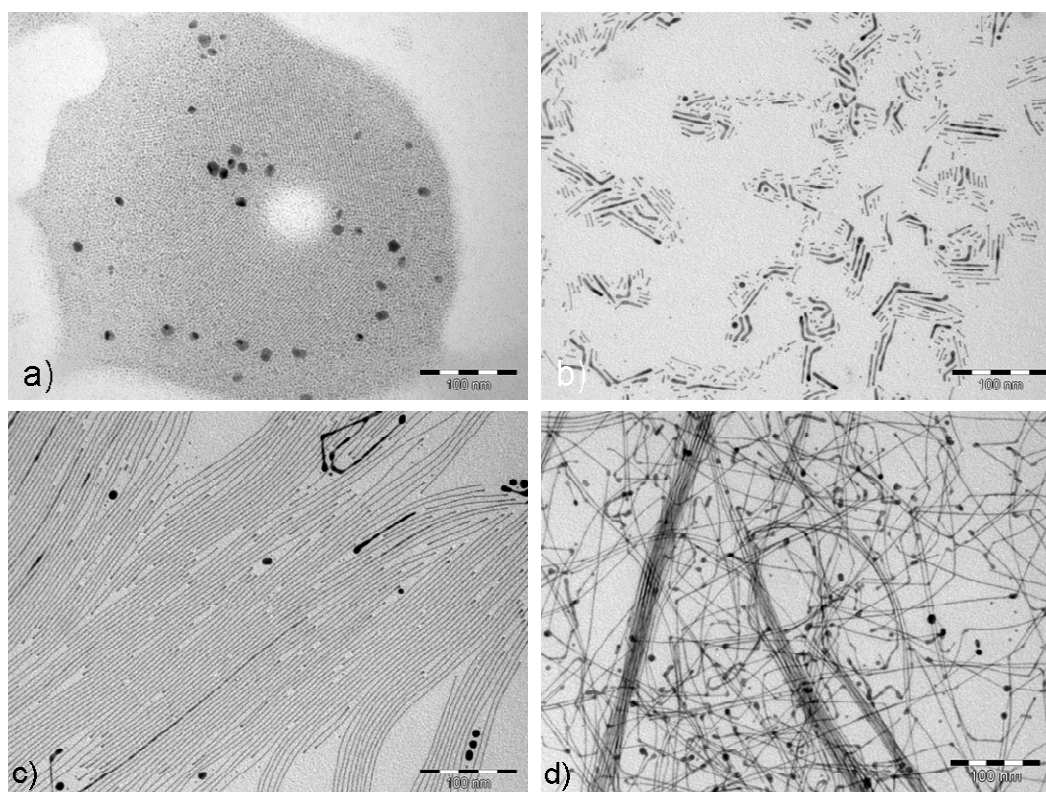


Figure S2: TEM image of Au NPs obtained after 5h at 40°C for TIPS/Au = 100 and [OY] = a) 20 ; b) 50 ; c) 100 and d) 250 mM.

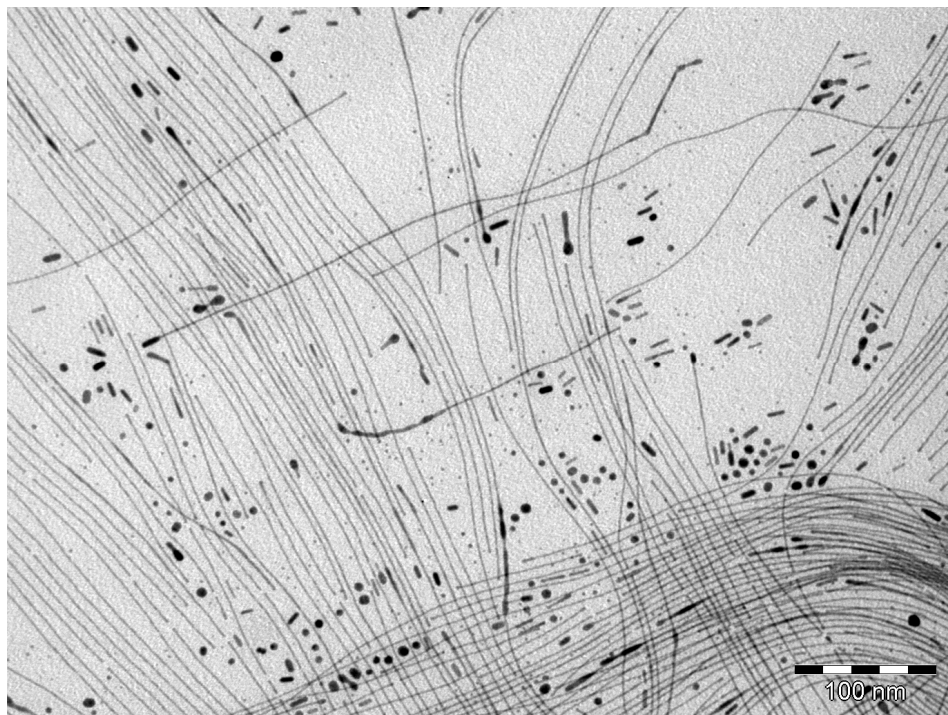


Figure S3: TEM image of Au NPs obtained after 27h at 40°C for TIPS/Au = 100 and [OY] = 400 mM.

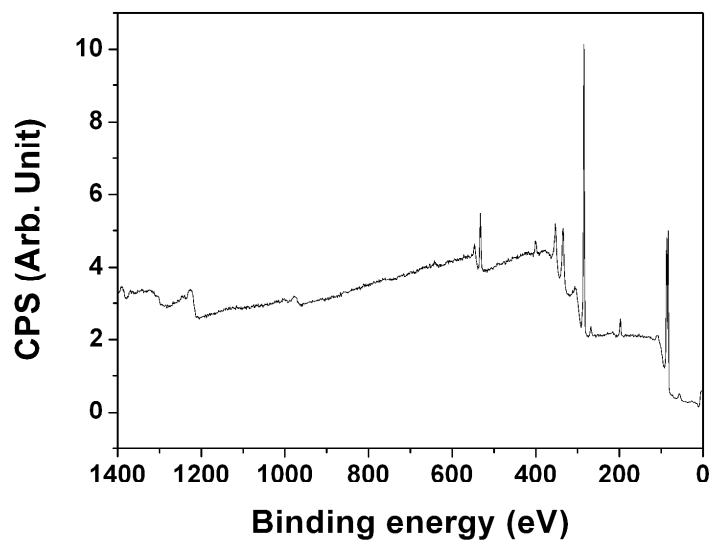


Figure S4: X-ray photon electron survey scan of Au nanowires precipitated by addition of absolute ethanol

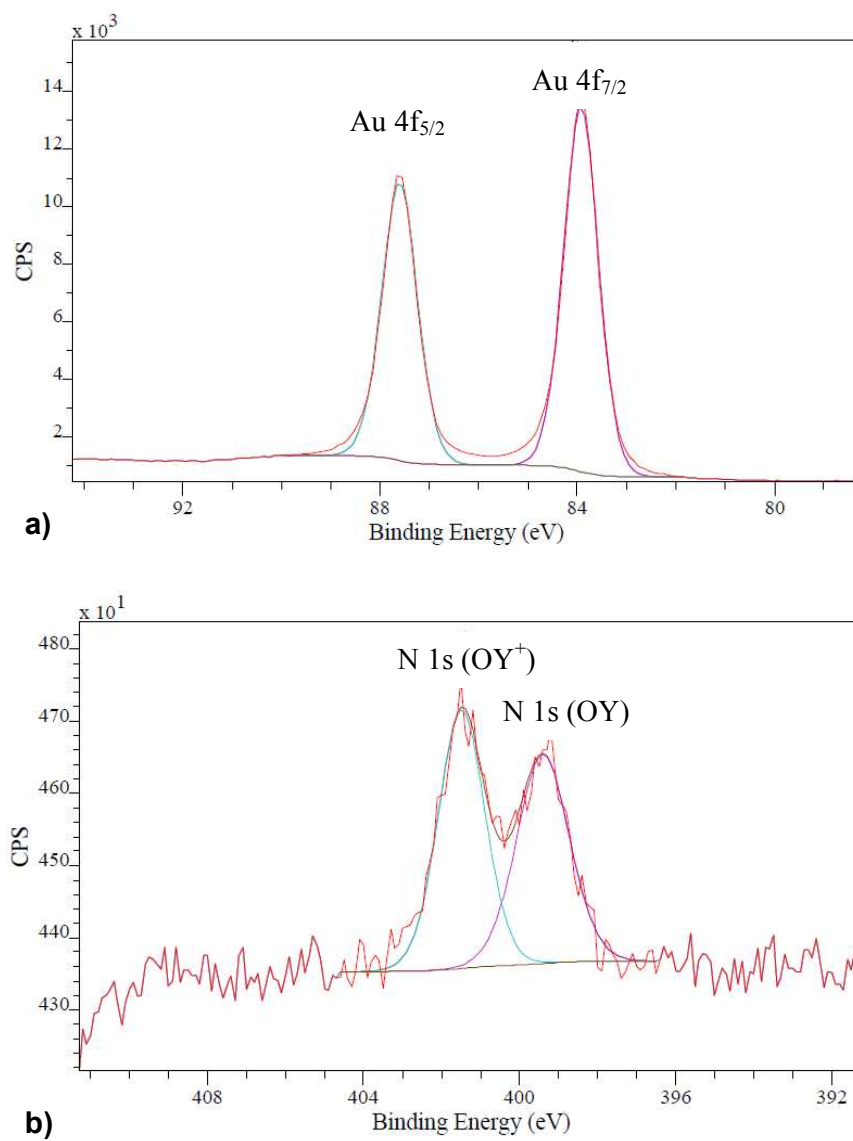


Figure S5: High resolution X-ray photon electron spectra of Au nanowires precipitated by addition of absolute ethanol (a) Au 4f peaks ; (b) N 1s peak and the corresponding fitting profiles

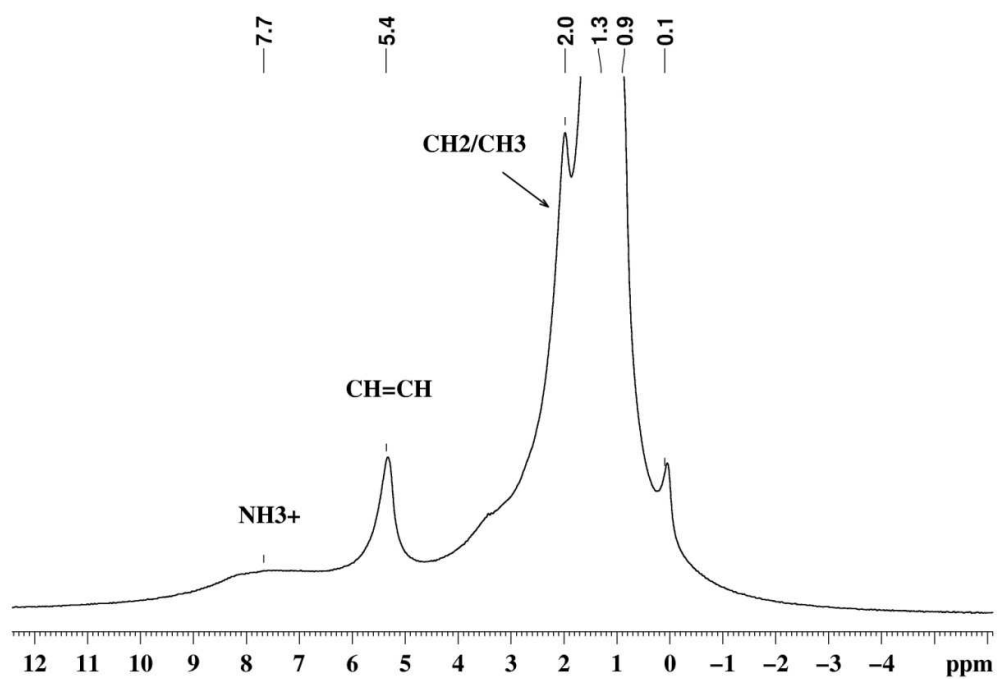


Figure S6: ^1H MAS spectrum (spinning rate 24 kHz) of Au NWs after purification.

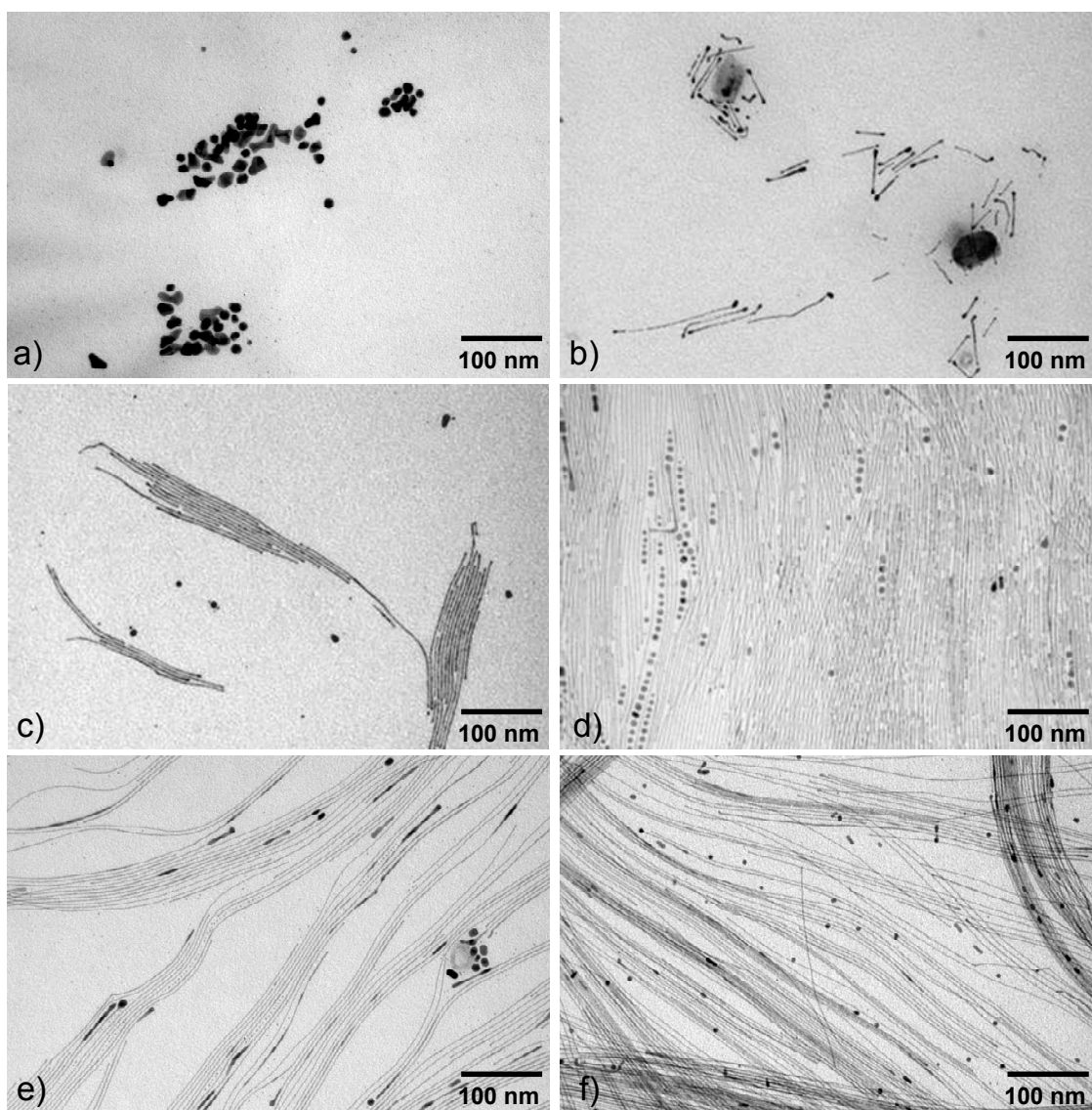
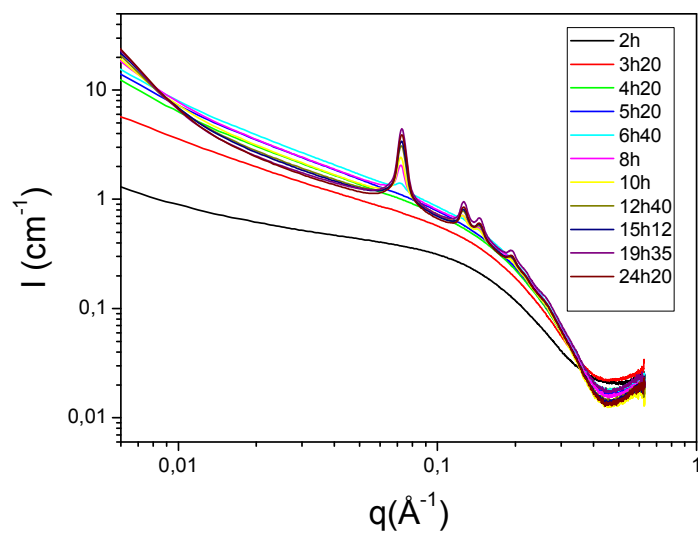
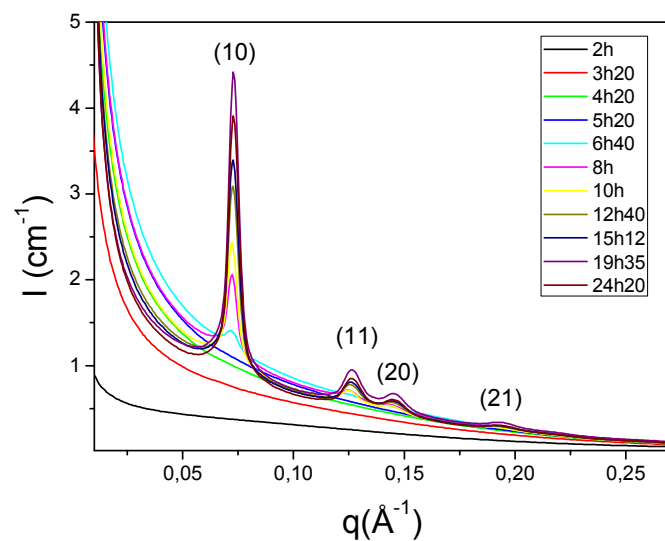


Figure S7. TEM images of Au nanoparticles obtained after a) 1, b) 2, c) 3, d) 4, e) 8 and f) 24 hours of reaction at 25 °C.



a)



b)

Figure S8: Small angle X-ray scattering results at 25 °C: (a) log-log scale ; (b) linear scale. The hexagonal phase appeared after 6h.

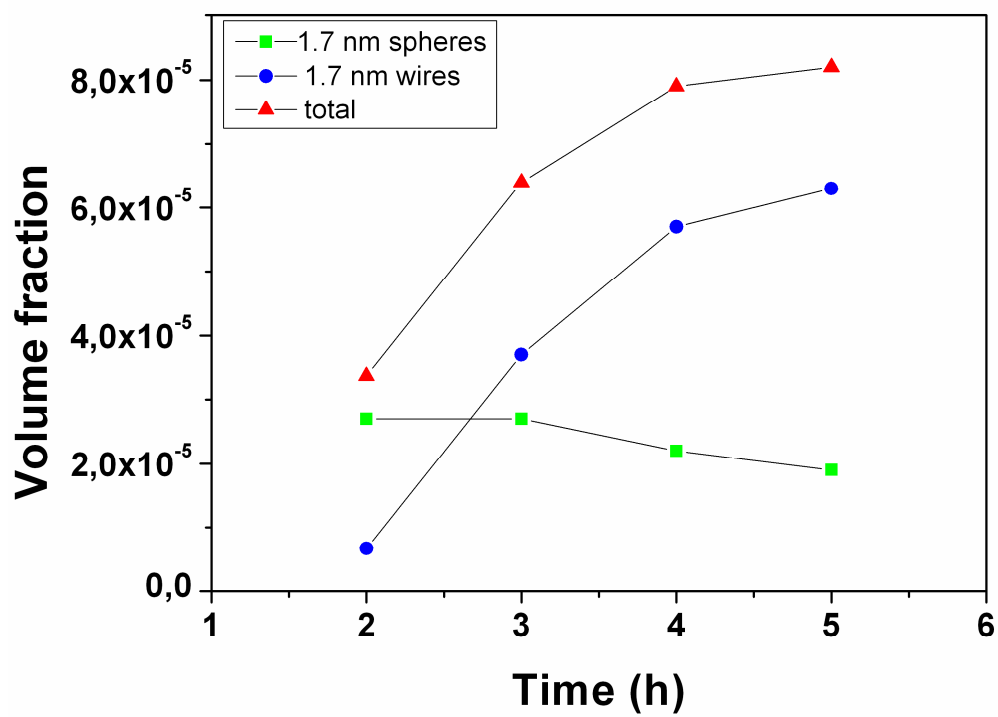


Figure S9: Volume fraction calculated from the SAXS modelling at different time of reaction at 25 °C.

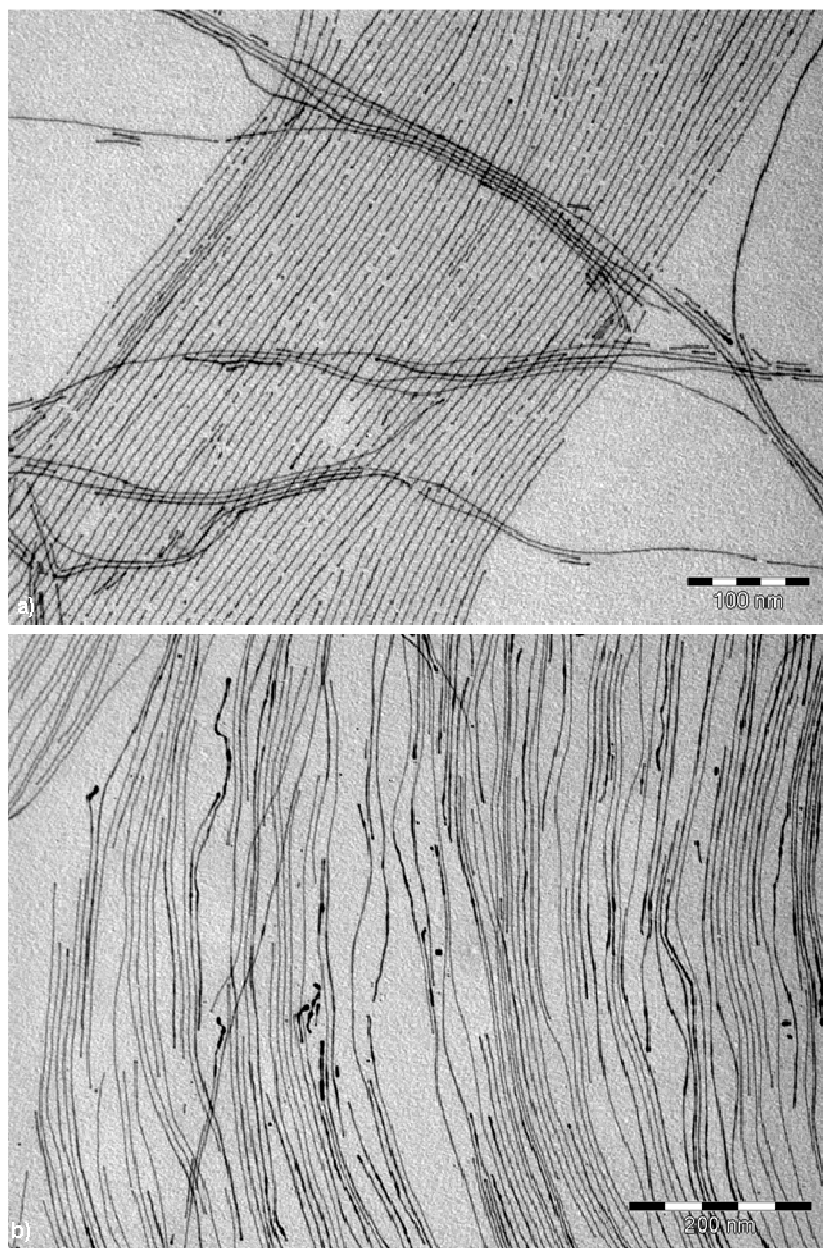


Figure S10. a-b) TEM image of assembled Au NWs.