

# Interstellar Chemistry: A Strategy for Detecting Polycyclic Aromatic Hydrocarbons in Space

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## Centrifugal Distortion Constants

The statistical weights for the  $K$  states of corannulene were determined using the analysis of Weber.<sup>1</sup> Corannulene ( $C_{5v}$ ) has atoms of type 5 in Table I of Weber.<sup>1</sup> Using Eq. (3) with  $h = 10$  (order of the group), Eqs. (4), (8), and (9) afford values for  $\chi_s(R)$ , and group theory analysis affords values for  $\chi_G^\alpha(R)$ .<sup>2</sup> The resulting  $m_\alpha$  values are:  $m(A_1) = 120$ ;  $m(A_2) = 88$ ;  $m(E_1) = m(E_2) = 204$ . Then from the first column of Table V<sup>1</sup> for the  $A_1$  ro-vibronic ground state:

$K$	<i>weight</i>
0	$m(A_1) + m(A_2) = 208$
$5p = 5, 10, \dots$	$2[m(A_1) + m(A_2)] = 416$
$5p \pm 1 = 1, 4, 6, 9, 11, \dots$	$2m(E_1) = 408$
$5p \pm 2 = 2, 3, 7, 8, 12, \dots$	$2m(E_2) = 408$

Thus, the weighting of the  $K = 0$  state is about half of all the other  $K$  states, and the relative magnitudes of the transition matrix elements follow the formula  $1 - [K^2 / (J + 1)^2]$ .

The analysis reveals that corannulene displays three symmetry species ( $A_1 + A_2$ ),  $E_1$ , and  $E_2$ , and cooling in the beam expansion might populate only the states available with energies up to the order of  $kT$ . In the case that  $D_{JK}$  is extremely small, the analysis is straightforward and the determined constants are  $B$  and  $D_J$ . In the case that  $D_{JK}$  is not small, the proper determination of  $D_J$  and  $D_{JK}$  becomes problematic. If  $D_{JK}$  is not small, the value determined for  $D_J$  is just an effective value, since the shift attributed to  $D_J$  might be caused by the population of higher  $K$  states (and depopulation of lower  $K$  states) as they become available with higher  $J$ . In the extreme case, only the three highest  $K$  states are populated after the transfer during the supersonic expansion. (If  $D_{JK}$  is sufficiently small, however, one would still see only one line in any  $J \rightarrow J + 1$  transition). If the population transfer is not complete, more  $K$  states would be

populated, but the “mean”  $K$ -value would not be zero. Most likely,  $K$  will be a function of  $J$ , e.g.

$K = J - X$  with  $J^2 > (J - X)^2 > \frac{kT}{C-B} + J^2 > 0$ , since states with a relative energy beyond the order of  $kT$  will essentially not be populated (e.g. for  $J$  levels with  $J^2 > \frac{kT}{B-C}$  the population of the  $K = 0$  state will decrease significantly; i.e.  $J > 14$  assuming a beam temperature of 2.5 K). Irrespective of whether or not population transfer is complete, a simple interpretation of this situation in terms of  $D_J$  and  $D_{JK}$  is not possible. In such a case, the value of  $D_J$  is just as speculative as  $D_{JK}$ , since it is not easily possible to distinguish in a fit between  $D_J \times J^2(J+1)^2$  and  $D_{JK} \times J(J+1)K^2$  approximately equal to  $D_{JK} \times J(J+1)(J-X)^2$ .

The Stark effect measurements, however, permit a definite statement on  $D_J$ . From the Stark measurements, the  $K = 0$  line position was extrapolated to zero-field conditions. Indeed, the extrapolated line coincides with the measured zero-field line position (Fig 4). This indicates the absence of a line shift arising from a mixture of different overlapping  $K$ -states, which confirms the small magnitude of  $D_{JK}$ . We therefore conclude that the line position for higher  $J$ -transitions is not affected by a mix of overlapping  $K$ -states.

From the line width at half maximum intensity for the highest frequency transition observed, a value of 0.0023 kHz for  $D_{JK}$  is determined as an upper limit to the splitting from  $K = 0$  to  $K = 18$ . This limit is half of the value of  $D_J$ . The upper limit for  $D_{JK}$  should probably be increased somewhat, since the highest  $K$  values are likely too weak to observe in the laboratory spectrum. In the end, extrapolation of the transition frequencies to the mm range (e.g. 75-100 GHz), as well as conclusions from the mm survey data, should be considered with due caution.

## Corannulene Bowl Inversion Process

Inversion along the c-axis (dipolar axis) would provide two levels for each J state.

Transitions from J" to J' would go from the bottom level in the lower state (J") to the upper level in the upper state (J'), and vice versa for the other two levels. Thus, the transitions go across the inversion splitting and the observed frequency splitting (if it happened) would be twice the inversion frequency. The narrowest linewidth measured in the current study is 9 kHz, which places an upper limit to the inversion frequency of 5 kHz at the temperature of the measurement (< 2 K). (The rotational temperature in the supersonic expansion is estimated to be 1-2 K).

From comparison with other systems, the low inversion frequency (< 5 kHz) seems reasonable for a large reduced mass tunneling through a barrier of 10.2 kcal/mol (measured for an analog of corannulene by dynamic NMR spectroscopy.<sup>3</sup>) Extrapolation of the measured rate to temperatures below 100 K predicts a vanishingly low rate of inversion via a classical, thermally-activated process. The absence of inversion splittings in the laboratory rotational spectra has no implication regarding astronomical observations, particularly for hot sources, which have broad lines. Dark clouds have narrower lines, which are, nevertheless, comparable to the 9 kHz width observed in Hannover.

## References

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**Table S1.** Calculated and Measured Rotational Transitions of Corannulene ( $C_{20}H_{10}$ ).

$J' - J''$	Calculated Frequency [MHz]	uncertainty (2 $\sigma$ )	Measured Frequency [MHz]	uncertainty (1 $\sigma$ )	Energy of Lower State [cm $^{-1}$ ]
1 - 0	1019.6854	( 1)			0.000
2 - 1	2039.3706	( 1)			0.034
3 - 2	3059.0556	( 2)			0.102
4 - 3	4078.7404	( 2)			0.204
5 - 4	5098.4247	( 2)	5098.4246	( 5)	0.340
6 - 5	6118.1084	( 3)	6118.1081	( 5)	0.510
7 - 6	7137.7916	( 3)	7137.7916	( 5)	0.714
8 - 7	8157.4740	( 3)	8157.4739	( 5)	0.952
9 - 8	9177.1556	( 3)	9177.156	( 2)	1.224
10 - 9	10196.8363	( 4)	10196.838	( 2)	1.531
11 - 10	11216.5159	( 4)	11216.517	( 2)	1.871
12 - 11	12236.1943	( 4)	12236.197	( 2)	2.245
13 - 12	13255.8715	( 5)	13255.872	( 2)	2.653
14 - 13	14275.5473	( 6)	14275.549	( 2)	3.095
15 - 14	15295.2217	( 7)	15295.222	( 2)	3.571
16 - 15	16314.8945	( 8)	16314.894	( 2)	4.082
17 - 16	17334.5657	( 10)	17334.564	( 2)	4.626
18 - 17	18354.2350	( 12)	18354.234	( 2)	5.204
19 - 18	19373.9025	( 15)	19373.903	( 2)	5.816
20 - 19	20393.5680	( 18)			6.462
21 - 20	21413.2314	( 21)			7.143
22 - 21	22432.8926	( 25)			7.857
23 - 22	23452.5515	( 29)			8.605
24 - 23	24472.2080	( 33)			9.388
25 - 24	25491.8619	( 39)			10.204
26 - 25	26511.5133	( 44)			11.054
27 - 26	27531.1620	( 50)			11.939
28 - 27	28550.8078	( 57)			12.857
29 - 28	29570.4507	( 64)			13.809
30 - 29	30590.0906	( 71)			14.796
31 - 30	31609.7273	( 79)			15.816
32 - 31	32629.3608	( 88)			16.870
33 - 32	33648.9910	( 97)			17.959
34 - 33	34668.6177	( 107)			19.081
35 - 34	35688.2408	( 118)			20.238
36 - 35	36707.8603	( 129)			21.428
37 - 36	37727.4760	( 141)			22.652
38 - 37	38747.0879	( 153)			23.911
39 - 38	39766.6957	( 166)			25.203
40 - 39	40786.2995	( 180)			26.530

41 - 40	41805.8992	( 195)	27.890
42 - 41	42825.4945	( 211)	29.285
43 - 42	43845.0854	( 227)	30.713
44 - 43	44864.6719	( 244)	32.176
45 - 44	45884.2537	( 262)	33.672
46 - 45	46903.8309	( 280)	35.203
47 - 46	47923.4032	( 300)	36.767
48 - 47	48942.9706	( 320)	38.366
49 - 48	49962.5330	( 341)	39.999
50 - 49	50982.0903	( 363)	41.665
51 - 50	52001.6424	( 387)	43.366
52 - 51	53021.1891	( 411)	45.100
53 - 52	54040.7304	( 436)	46.869
54 - 53	55060.2661	( 461)	48.671
55 - 54	56079.7962	( 488)	50.508
56 - 55	57099.3205	( 516)	52.379
57 - 56	58118.8390	( 545)	54.283
58 - 57	59138.3516	( 575)	56.222
59 - 58	60157.8580	( 607)	58.195
60 - 59	61177.3583	( 639)	60.201
61 - 60	62196.8524	( 672)	62.242
62 - 61	63216.3400	( 706)	64.317
63 - 62	64235.8212	( 742)	66.425
64 - 63	65255.2957	( 779)	68.568
65 - 64	66274.7636	( 817)	70.745
66 - 65	67294.2247	( 856)	72.955
67 - 66	68313.6789	( 896)	75.200
68 - 67	69333.1261	( 938)	77.479
69 - 68	70352.5662	( 981)	79.791
70 - 69	71371.9991	( 1025)	82.138
71 - 70	72391.4246	( 1070)	84.519
72 - 71	73410.8427	( 1117)	86.934
73 - 72	74430.2533	( 1165)	89.382
74 - 73	75449.6563	( 1214)	91.865
75 - 74	76469.0515	( 1265)	94.382
76 - 75	77488.4389	( 1317)	96.932
77 - 76	78507.8183	( 1370)	99.517
78 - 77	79527.1897	( 1425)	102.136
79 - 78	80546.5530	( 1482)	104.789
80 - 79	81565.9079	( 1540)	107.475
81 - 80	82585.2545	( 1599)	110.196
82 - 81	83604.5927	( 1660)	112.951
83 - 82	84623.9223	( 1722)	115.740
84 - 83	85643.2432	( 1786)	118.562
85 - 84	86662.5553	( 1851)	121.419
86 - 85	87681.8585	( 1918)	124.310
87 - 86	88701.1527	( 1987)	127.235
88 - 87	89720.4379	( 2057)	130.193
89 - 88	90739.7138	( 2129)	133.186
90 - 89	91758.9804	( 2202)	136.213

91 - 90	92778.2376	( 2277)	139.274
92 - 91	93797.4854	( 2354)	142.368
93 - 92	94816.7234	( 2432)	145.497
94 - 93	95835.9518	( 2512)	148.660
95 - 94	96855.1703	( 2594)	151.857
96 - 95	97874.3790	( 2678)	155.087
97 - 96	98893.5775	( 2763)	158.352
98 - 97	99912.7660	( 2850)	161.651
99 - 98	100931.9441	( 2939)	164.984
100 - 99	101951.1120	( 3030)	168.350
101 - 100	102970.2693	( 3123)	171.751
102 - 101	103989.4162	( 3217)	175.186
103 - 102	105008.5523	( 3314)	178.654
104 - 103	106027.6777	( 3412)	182.157
105 - 104	107046.7922	( 3512)	185.694
106 - 105	108065.8957	( 3614)	189.265
107 - 106	109084.9882	( 3719)	192.869
108 - 107	110104.0695	( 3825)	196.508
109 - 108	111123.1394	( 3933)	200.181
110 - 109	112142.1980	( 4043)	203.887
111 - 110	113161.2451	( 4155)	207.628
112 - 111	114180.2806	( 4269)	211.403
113 - 112	115199.3043	( 4385)	215.211
114 - 113	116218.3163	( 4504)	219.054
115 - 114	117237.3163	( 4624)	222.931
116 - 115	118256.3043	( 4746)	226.841
117 - 116	119275.2802	( 4871)	230.786
118 - 117	120294.2438	( 4998)	234.764
119 - 118	121313.1952	( 5127)	238.777
120 - 119	122332.1340	( 5258)	242.823
121 - 120	123351.0604	( 5392)	246.904
122 - 121	124369.9740	( 5527)	251.019
123 - 122	125388.8749	( 5665)	255.167
124 - 123	126407.7630	( 5805)	259.350
125 - 124	127426.6381	( 5948)	263.566
126 - 125	128445.5001	( 6092)	267.817
127 - 126	129464.3490	( 6239)	272.101
128 - 127	130483.1846	( 6389)	276.420
129 - 128	131502.0067	( 6541)	280.772
130 - 129	132520.8155	( 6695)	285.158
131 - 130	133539.6106	( 6851)	289.579
132 - 131	134558.3920	( 7010)	294.033
133 - 132	135577.1596	( 7171)	298.522
134 - 133	136595.9133	( 7335)	303.044
135 - 134	137614.6530	( 7502)	307.600
136 - 135	138633.3786	( 7670)	312.191
137 - 136	139652.0900	( 7842)	316.815
138 - 137	140670.7870	( 8016)	321.473
139 - 138	141689.4696	( 8192)	326.166
140 - 139	142708.1377	( 8371)	330.892

141 - 140	143726.7912	( 8552)	335.652
142 - 141	144745.4299	( 8736)	340.446
143 - 142	145764.0538	( 8923)	345.275
144 - 143	146782.6627	( 9112)	350.137
145 - 144	147801.2565	( 9304)	355.033
146 - 145	148819.8352	( 9499)	359.963
147 - 146	149838.3987	( 9697)	364.927
148 - 147	150856.9467	( 9897)	369.925
149 - 148	151875.4793	( 10099)	374.957
150 - 149	152893.9963	( 10305)	380.023
151 - 150	153912.4977	( 10513)	385.123
152 - 151	154930.9832	( 10724)	390.257
153 - 152	155949.4529	( 10938)	395.425
154 - 153	156967.9065	( 11155)	400.627
155 - 154	157986.3441	( 11375)	405.863

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