

## Supplementary material for *Identifying the significance of nonlinear normal modes*

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### Conditionally-resonant terms and phase dependence

For an  $N$ -DOF system, the  $\ell^{\text{th}}$  element of the vector of variables,  $\mathbf{u}^*$ , may be written

$$u_\ell^* = \prod_{k=1}^N u_{pk}^{s_{p\ell,k}} u_{mk}^{s_{m\ell,k}}, \quad (0.1)$$

where  $u_{pk}$  and  $u_{mk}$  are defined as

$$u_{pk} = \frac{U_k}{2} \exp \{ +j (\omega_{rk} t - \phi_k) \}, \quad u_{mk} = \frac{U_k}{2} \exp \{ -j (\omega_{rk} t - \phi_k) \}, \quad (0.2)$$

where  $U_k$ ,  $\omega_{rk}$  and  $\phi_k$  represent the amplitude, response frequency and phase of the fundamental component of the  $k^{\text{th}}$  linear mode respectively.

Substituting Eqs. (0.2) into Eq. (0.1) gives

$$u_\ell^* = \left[ \prod_{k=1}^N \left( \frac{U_k}{2} \right)^{(s_{p\ell,k} + s_{m\ell,k})} \right] \exp \left\{ +j \sum_{k=1}^N \tilde{s}_{\ell,k} (\omega_{rk} t - \phi_k) \right\}, \quad (0.3)$$

where  $\tilde{s}_{\ell,k} = s_{p\ell,k} - s_{m\ell,k}$ . Using this, element  $\{i, \ell\}$  of  $\beta$  may be found using

$$\beta_{i,\ell} = \left( \sum_{k=1}^N \tilde{s}_{\ell,k} \omega_{rk} \right)^2 - \omega_{ri}^2, \quad (0.4)$$

where  $\beta_{i,\ell}$  represents a resonant term when  $\beta_{i,\ell} = 0$ . We now define  $r_k$  such that  $\omega_{rk} = r_k \omega$ , where  $\omega$  represents the base frequency i.e.  $\omega = 2\pi/T$  and where  $T$  is the period of the response. Substituting this into Eq. (0.4),  $\beta_{i,\ell}$  represents a resonant term when

$$\left( \sum_{k=1}^N \tilde{s}_{\ell,k} r_k \right)^2 = r_i^2. \quad (0.5)$$

If the term represented by this is unconditionally-resonant, then Eq. (0.5) must be satisfied regardless of the values of  $r_k$ . It can be seen that this is only true when

$$\tilde{s}_{\ell,k} = \begin{cases} 0 & \text{for } k \neq i, \\ \pm 1 & \text{for } k = i. \end{cases} \quad (0.6)$$

For a conservative system, the vector of resonant nonlinear terms,  $\mathbf{N}_u$ , is written

$$\mathbf{N}_u(\mathbf{u}) = [\mathbf{N}_u] \mathbf{u}^*, \quad (0.7)$$

where  $[\mathbf{N}_u]$  is an  $\{N \times L\}$  matrix of coefficients whose  $\{i, \ell\}^{\text{th}}$  element is written  $[\mathbf{N}_u]_{i,\ell}$ . Therefore the  $i^{\text{th}}$  element of  $\mathbf{N}_u$  is given by

$$N_{ui} = \sum_{\ell=1}^L [\mathbf{N}_u]_{i,\ell} u_\ell^*. \quad (0.8)$$

Using Eq. (0.3), this may be written

$$N_{ui} = \sum_{\ell=1}^L K_{i,\ell} \exp \left\{ +j \sum_{k=1}^N \tilde{s}_{\ell,k} (\omega_{rk} t - \phi_k) \right\}. \quad (0.9)$$

If all terms are unconditionally-resonant (i.e. there are no conditionally-resonant terms in the resonant equation of motion) then, using Eq. (0.6), Eq. (0.9) may be written

$$N_{ui} = \sum_{\ell=1}^L [N_u]_{i,\ell} \left[ \prod_{k=1}^N \left( \frac{U_k}{2} \right)^{(s_{p\ell,k} + s_{m\ell,k})} \right] e^{\pm j(\omega_{ri}t - \phi_i)}, \quad (0.10)$$

The  $i^{\text{th}}$  resonant equation of motion is written

$$\ddot{u}_i + \omega_{ni}^2 u_i + N_{ui} = 0, \quad (0.11)$$

which, substituting the assumed solution for  $u_i$  and Eq. (0.10), gives

$$\left\{ \left( \omega_{ni}^2 - \omega_{ri}^2 \right) \frac{U_i}{2} + \sum_{\ell=1}^L [N_u]_{i,\ell} \left[ \prod_{k=1}^N \left( \frac{U_k}{2} \right)^{(s_{p\ell,k} + s_{m\ell,k})} \right] \right\} e^{\pm j(\omega_{ri}t - \phi_i)} = 0. \quad (0.12)$$

The time-independent components of Eq. (0.12) may then be written

$$\left( \omega_{ni}^2 - \omega_{ri}^2 \right) \frac{U_i}{2} + \sum_{\ell=1}^L [N_u]_{i,\ell} \left[ \prod_{k=1}^N \left( \frac{U_k}{2} \right)^{(s_{p\ell,k} + s_{m\ell,k})} \right] = 0. \quad (0.13)$$

It is clear that Eq. (0.13) is independent of phase. This demonstrates that, if a resonant equation of motion contains only unconditionally-resonant terms (i.e. it does not contain any conditionally-resonant terms), the solution will be phase-unlocked. As such, it may be concluded that only conditionally-resonant terms may lead phase-locking terms, and hence conditionally-resonant terms are required for phase-locked NNMs.